

Moments of inertia of pairing rotation within the BCS model for Sn and Ni isotopes

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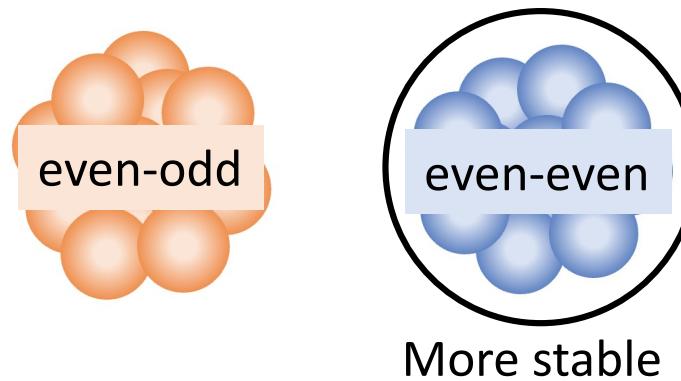
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Introduction

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Due to pairing interaction

More stable

Only even-even nuclei

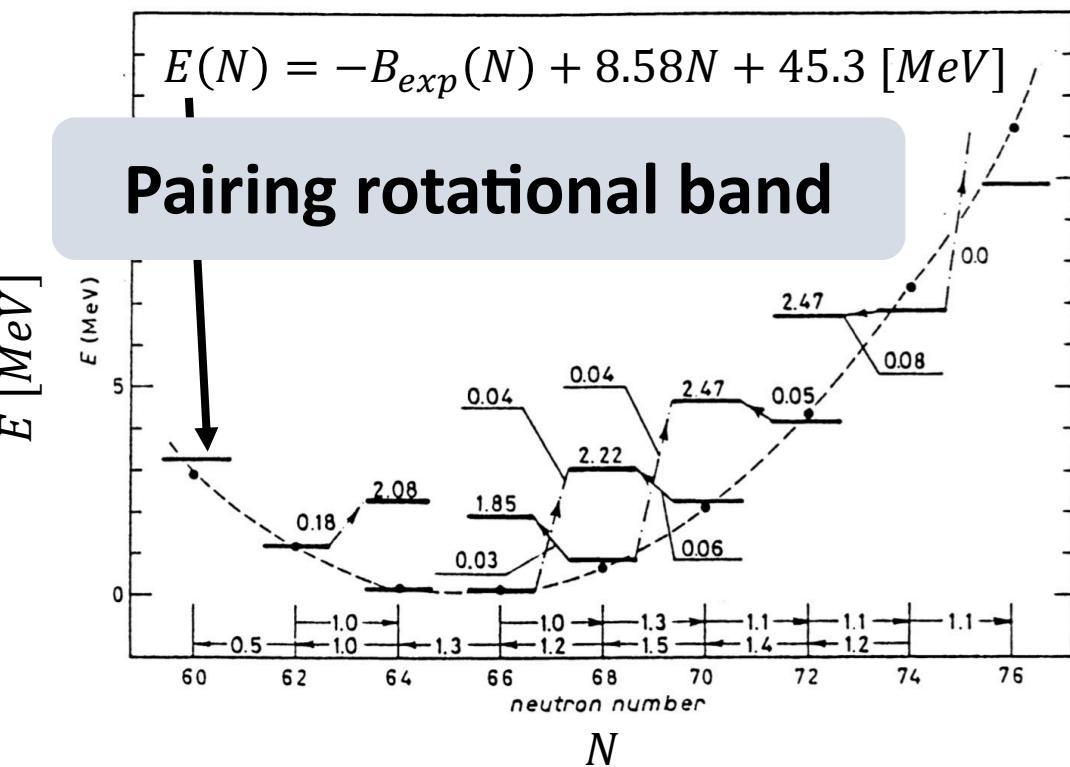
Ground state energy E_0

$$E_0(N) - E_0(N_0) - \frac{\partial E_0}{\partial N}(N - N_0)$$

$$= \frac{1}{2} \frac{\partial^2 E_0}{\partial N^2} (N - N_0)^2 \dots$$

Second-order term of E_0 w.r.t. N

Pairing rotational energy



Pairing rotational energy in Sn isotopes

Introduction

The pair condensation in a nucleus

\leftrightarrow “**deformation**” in the gauge space



Creating a new rotational degree of freedom in the gauge space



The energy of this rotation = **pairing rotational energy**

Energy of rotation
in real space

$$E(J) = \frac{J(J+1)}{2I}$$

J : angular momentum

I : moment of inertia

Energy of rotation
in gauge space

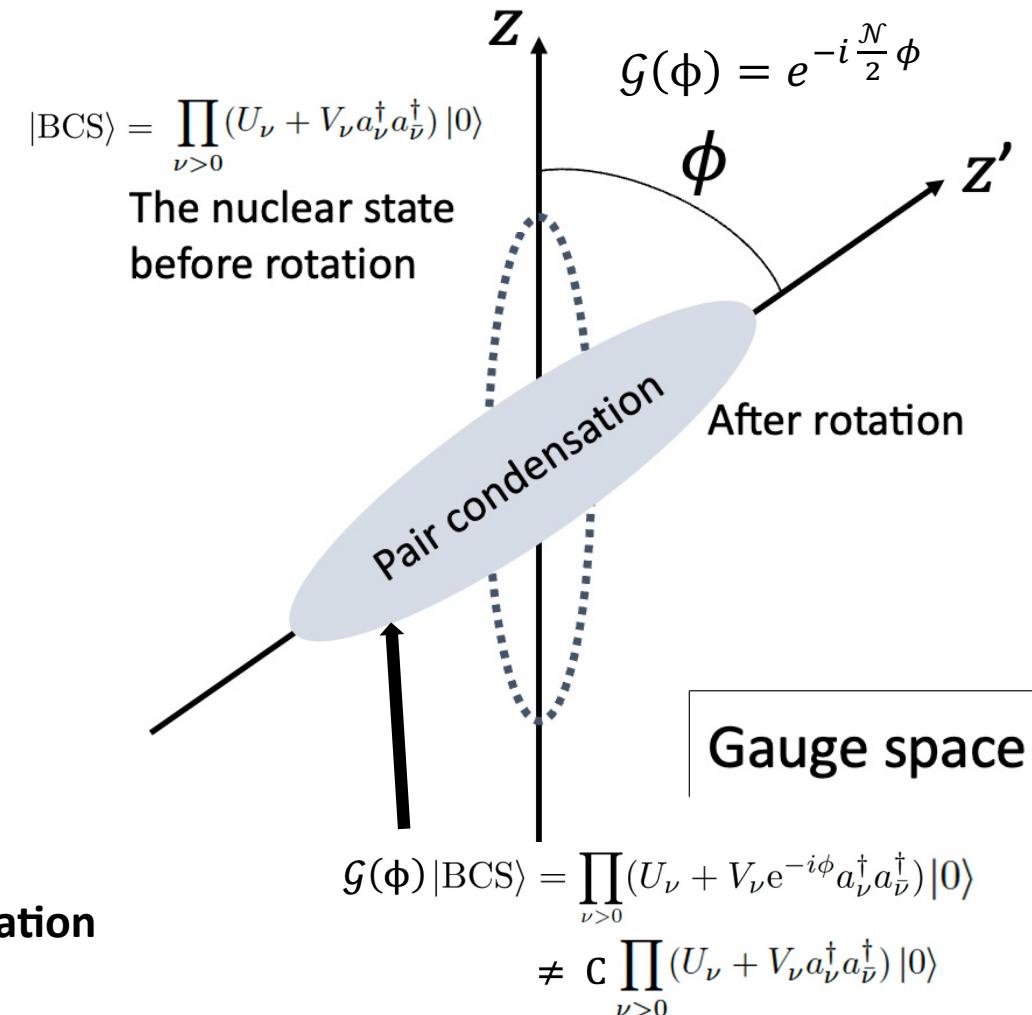
$$E(N) = \frac{(N - N_0)^2}{2J}$$

N : particle number

J : **Moment of inertia of pairing rotation**

$$E_0(N) = E_0(N_0) + \frac{\partial E_0}{\partial N} (N - N_0) + \frac{(N - N_0)^2}{2J}$$

$E(N)$ corresponds to
second-order term of $E_0(N)$

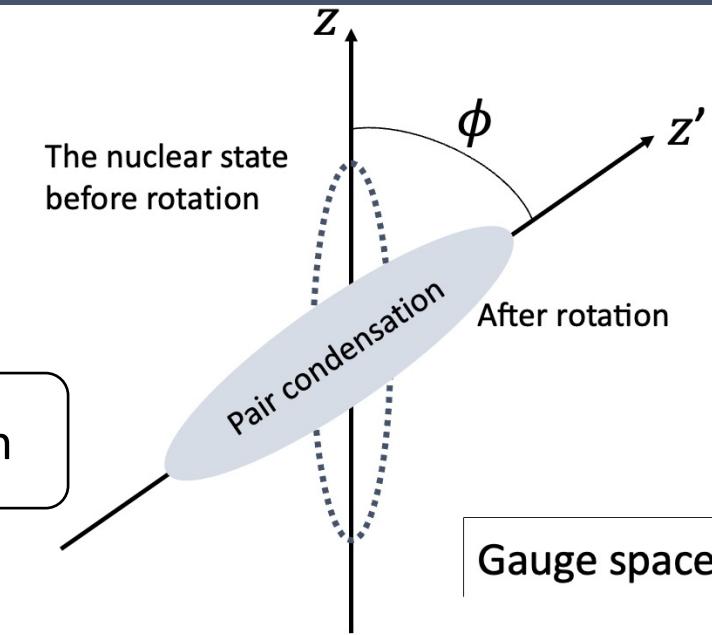


Purpose

$$E_0(N) = E_0(N_0) + \frac{\partial E_0}{\partial N} (N - N_0) + \frac{(N - N_0)^2}{2J}$$

Pairing rotational energy

This quantity has not received much attention



To Investigate the various properties of J
by comparing with those of spatial rotation

Method

Hamiltonian

$$H = \sum_{\nu>0} e_\nu (a_\nu^\dagger a_\nu + a_{\bar{\nu}}^\dagger a_{\bar{\nu}}) - G \sum_{\mu>0} \sum_{\nu>0} a_\mu^\dagger a_{\bar{\mu}}^\dagger a_{\bar{\nu}} a_\nu$$

Pairing term

BCS model

$$|\text{BCS}\rangle = \prod_{\nu>0} (U_\nu + V_\nu a_\nu^\dagger a_{\bar{\nu}}^\dagger) |0\rangle$$

→ Find the ground state energy from the variational principle

$$\delta \langle \text{BCS} | H - \lambda N | \text{BCS} \rangle = 0 \quad \langle \text{BCS} | N | \text{BCS} \rangle = n$$

↳ Condition of the number of particles

Gap Equation

$$\Delta = \frac{1}{2} G \sum_{\nu>0} \Omega_\nu \frac{\Delta}{\sqrt{(e_\nu - \lambda)^2 + \Delta^2}}$$

☞ Calculate self-consistently

where $\Delta = G \sum_{\nu>0} \Omega_\nu U_\nu V_\nu$

→ Determine U, V, and calculate ground state energy E_0

$$E_0 = 2 \sum_\nu \Omega_\nu e_\nu V_\nu^2 - \frac{\Delta^2}{G} - G \sum_\nu \Omega_\nu V_\nu^4$$

→ Calculate **MOI of pairing rotation**

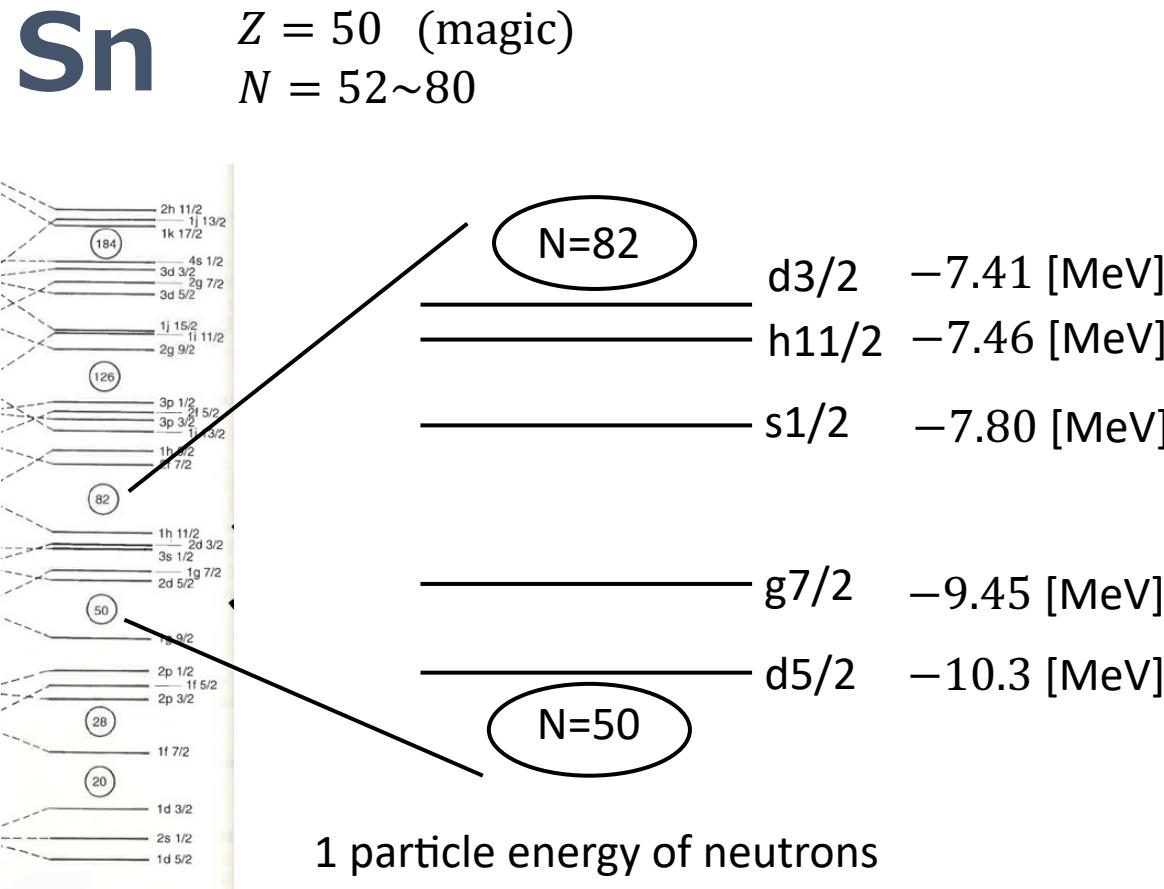
$$\mathcal{J} = \frac{1}{\frac{\partial^2 E_0}{\partial N^2}} = \frac{4}{E_0(N-2) - 2E_0(N) + E_0(N+2)}$$

$$\begin{aligned} \Delta S_{2n}(N) &= S_{2n}(N) - S_{2n}(N+2) \\ &= E_0(N-2) - 2E_0(N) + E_0(N+2) \end{aligned}$$

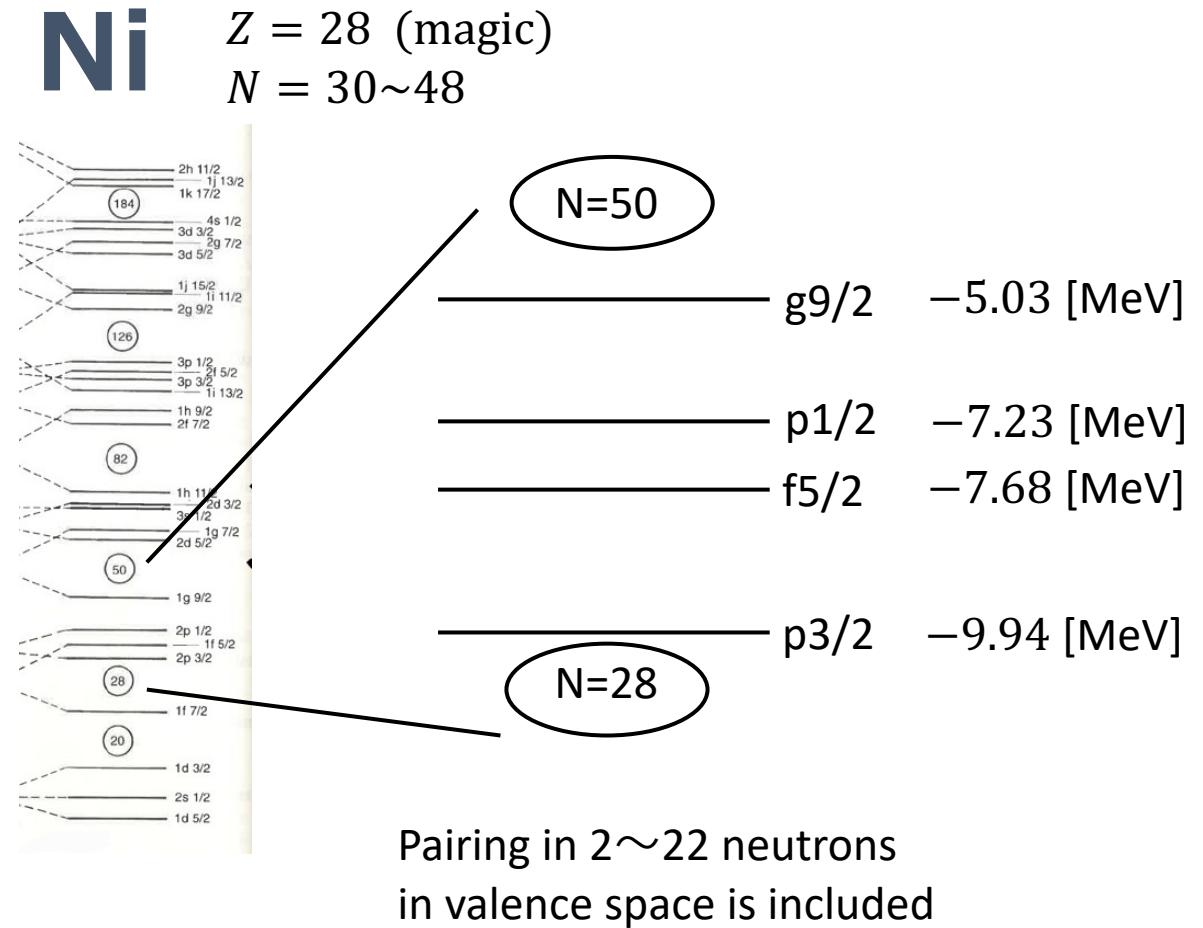
→ $\Delta S_{2n} = 4/\mathcal{J}$

Method

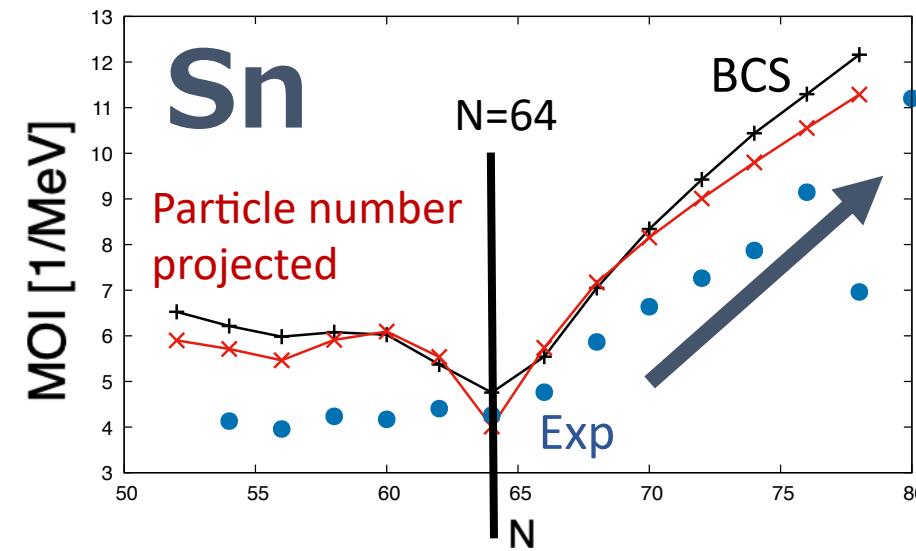
Only neutron pairing is considered



Pairing in $2 \sim 30$ neutrons
in valence space is included



N dependence of MOI



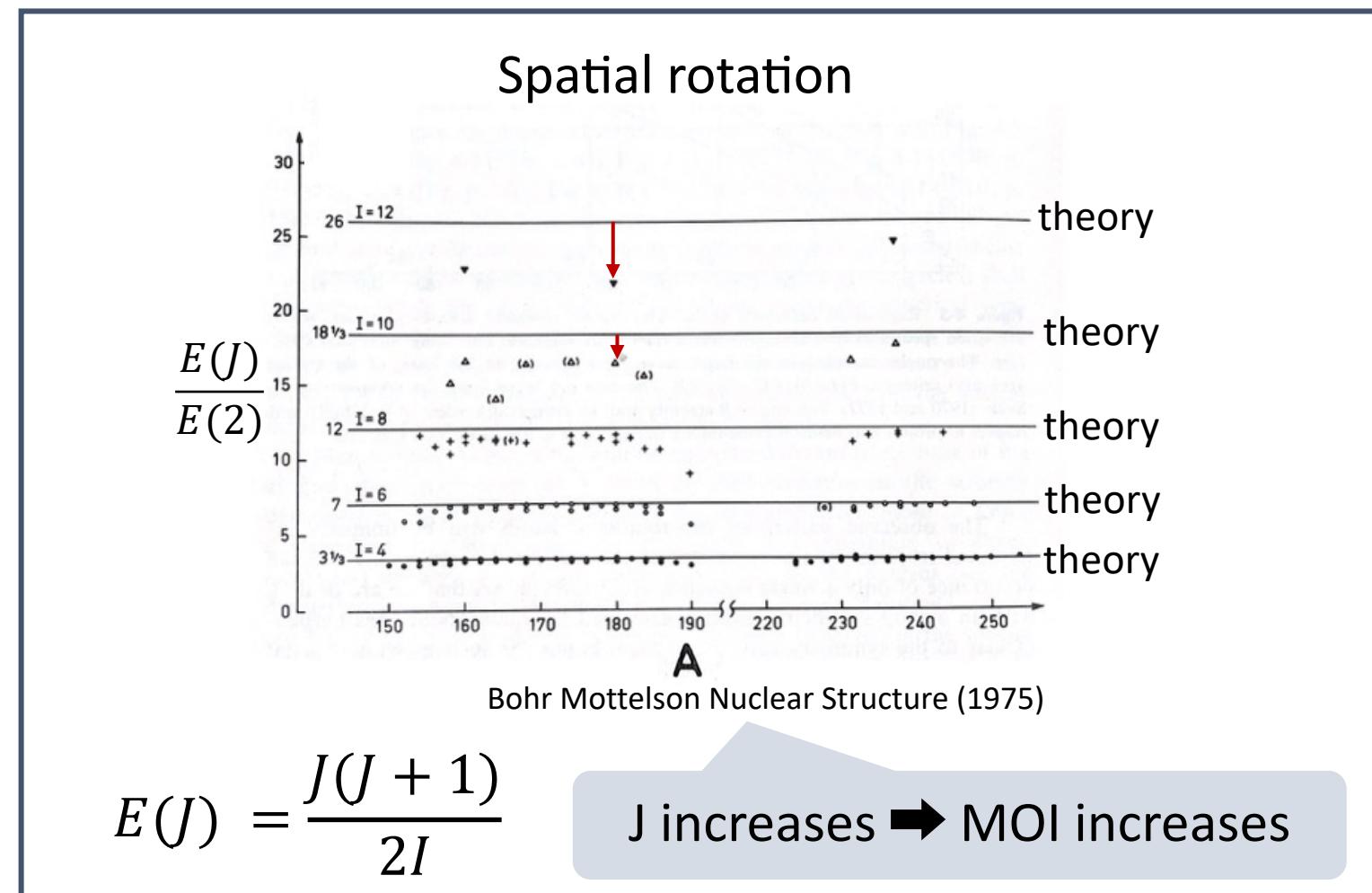
Particle number projected :

$$E_0(\text{projected}) = \langle n | H | n \rangle$$

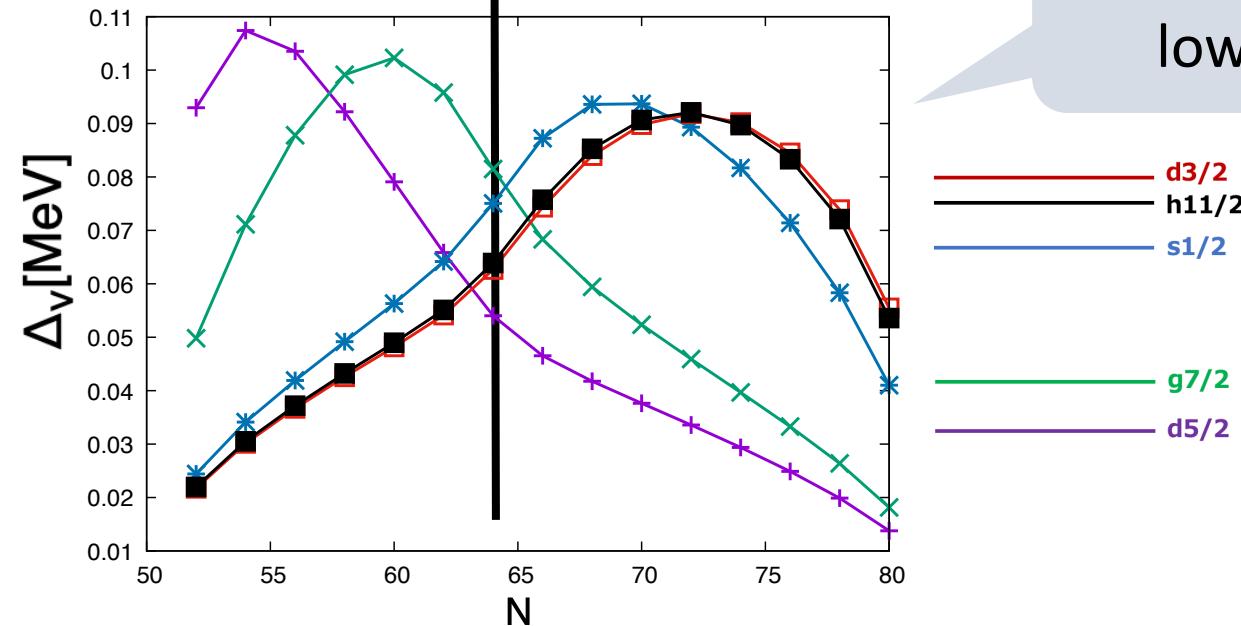
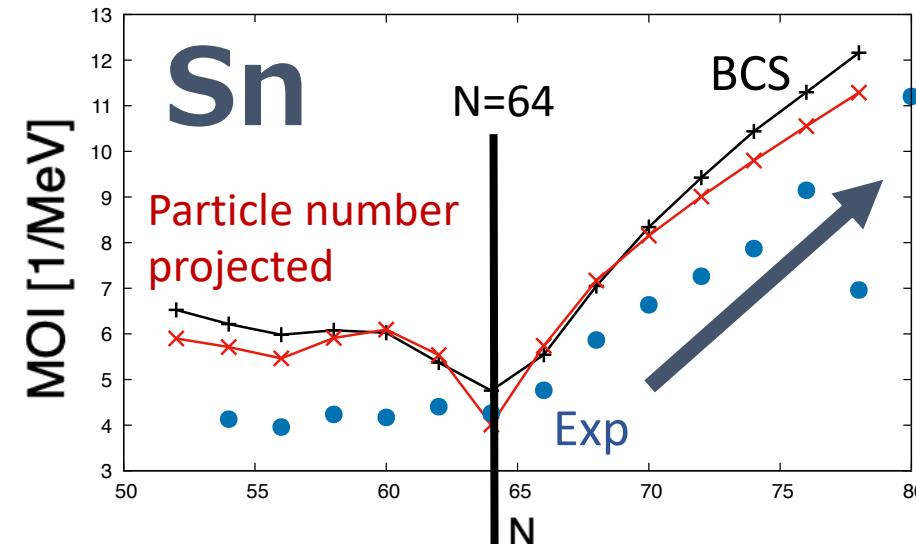
Particle number eigen state
(after projected)

N increases \rightarrow MOI increases

Same as the relationship between J -MOI in spatial



N dependence of MOI

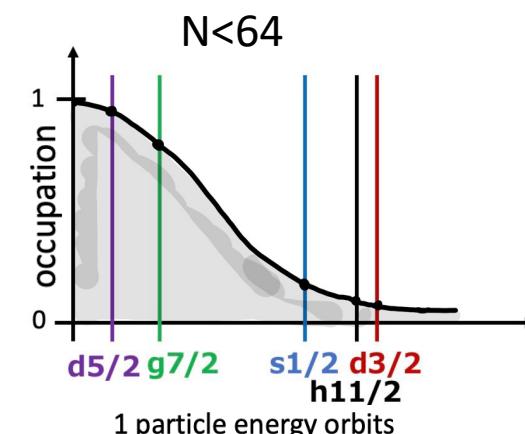


Above $N=64$, the increasing trend of MOI changes

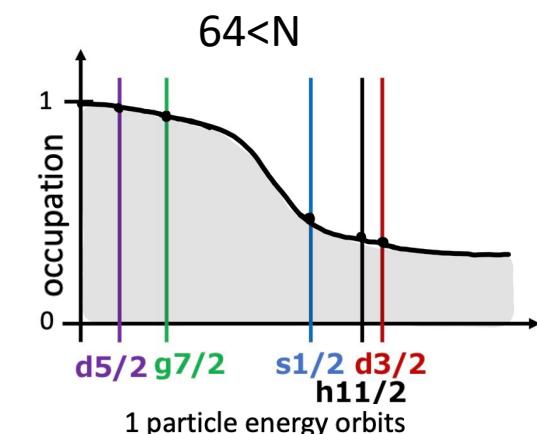
$$\Delta = \sum_{\nu>0} \Omega_\nu \Delta_\nu$$

Pairing gap of each orbit

The neutron contributing to superfluidity moves from lower 2 energy orbits to higher 3 energy orbits

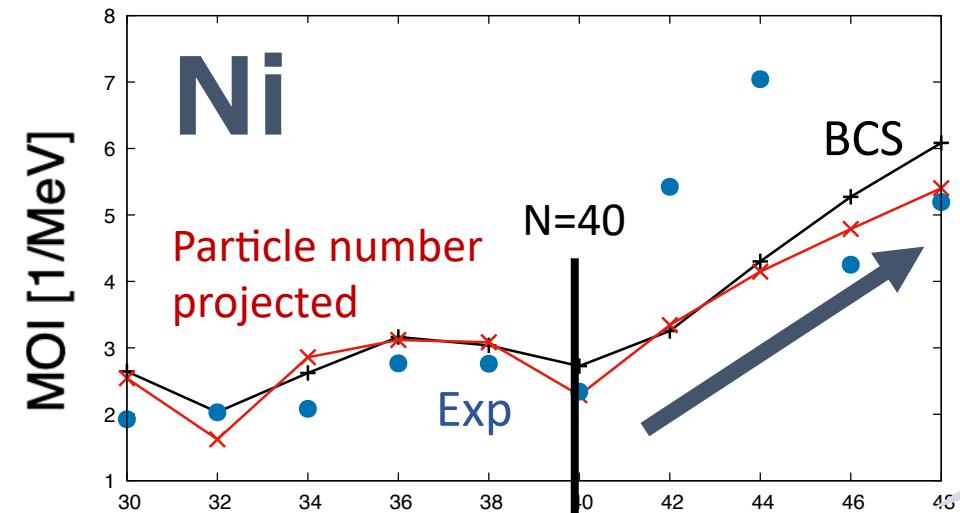


$N=64$



Why?

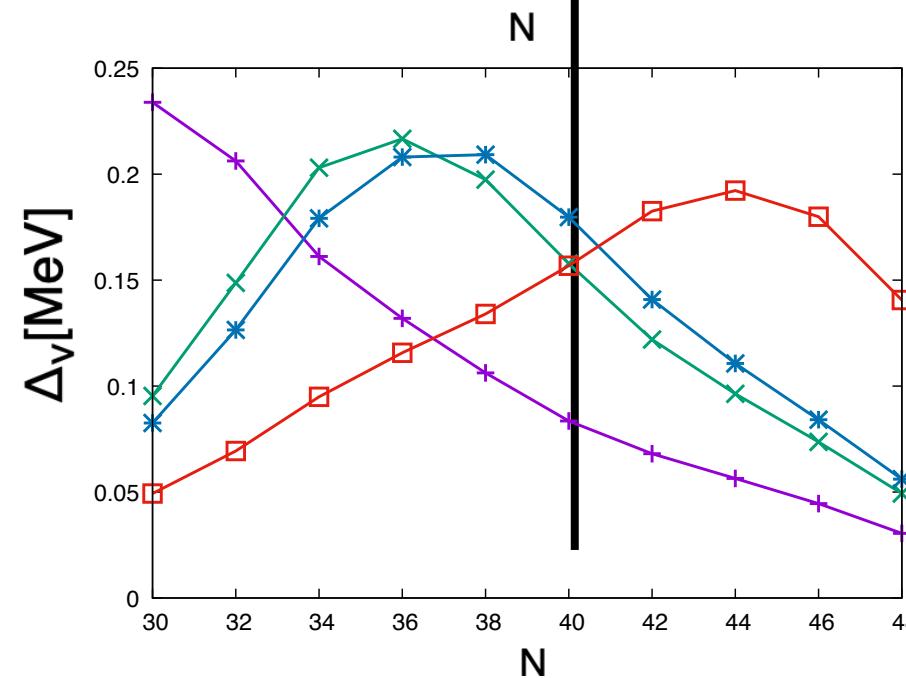
N dependence of MOI



- N increases → MOI increases
- After N=40, the increasing trend of MOI changes

Same as Sn !

The neutron contributing to superfluidity moves from lower 3 energy orbits to higher 1 energy orbits



$$\Delta = \sum_{\nu>0} \Omega_\nu \Delta_\nu$$

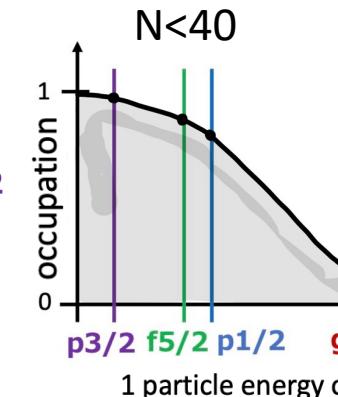
Pairing gap
of each orbit

g9/2

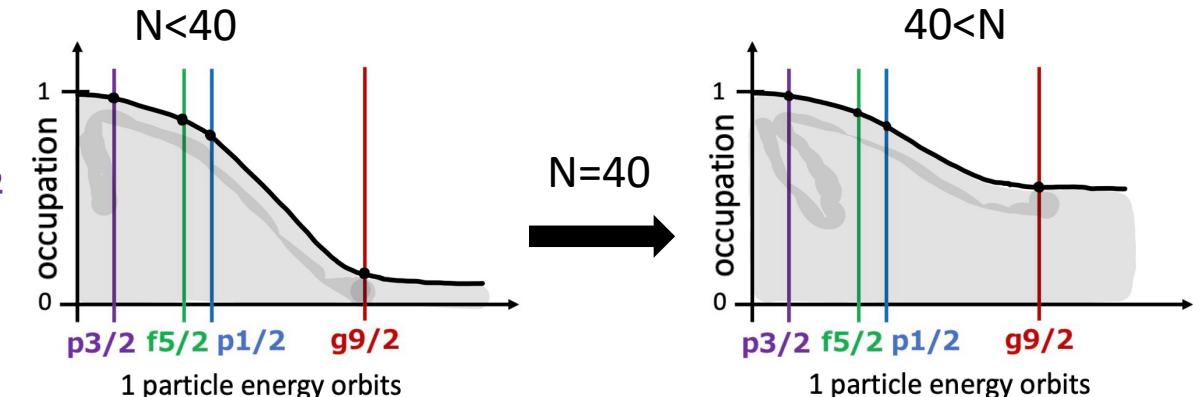
p1/2

f5/2

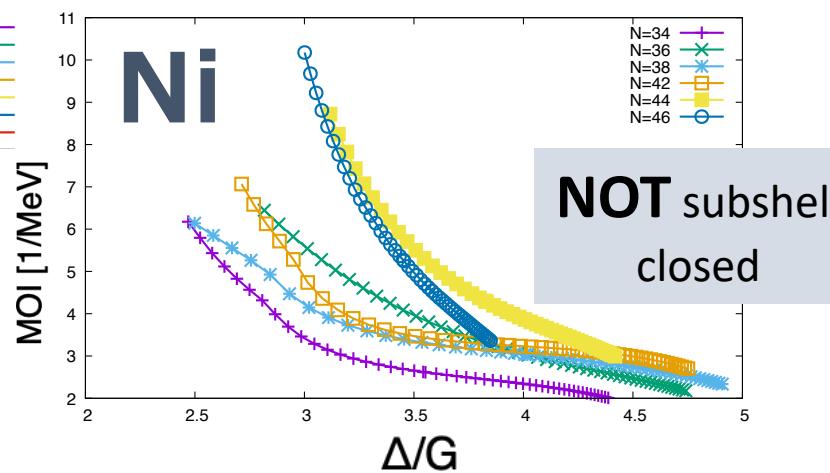
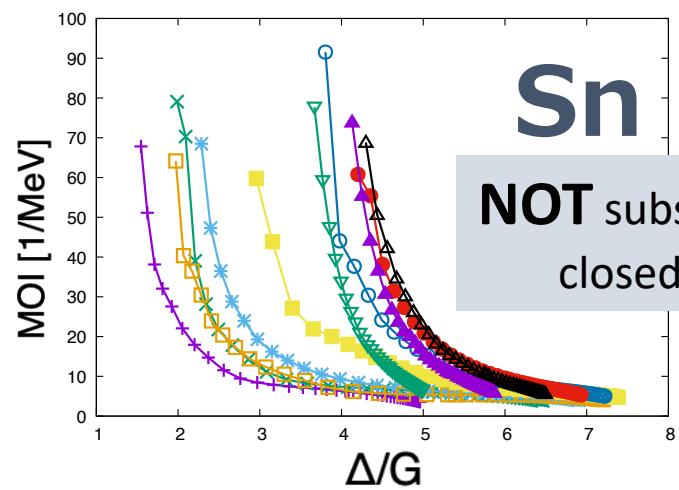
p3/2



$N=40$

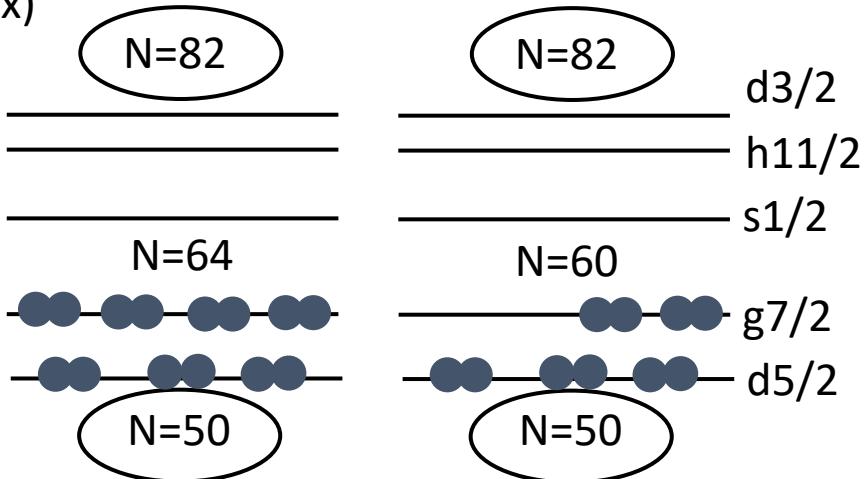


Δ dependence of MOI



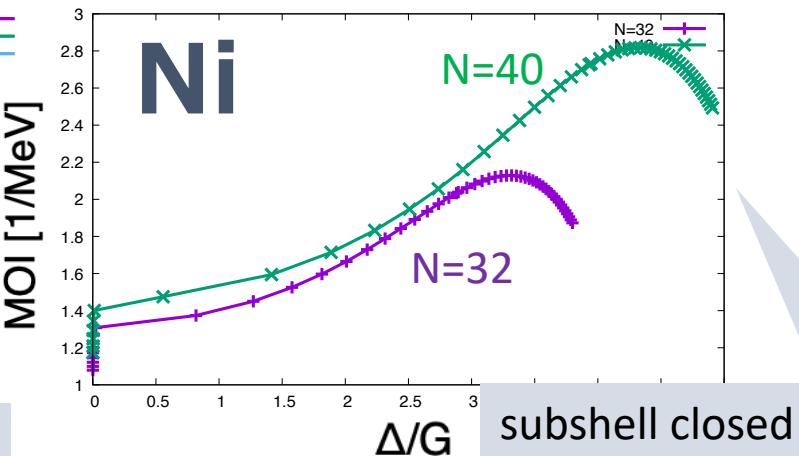
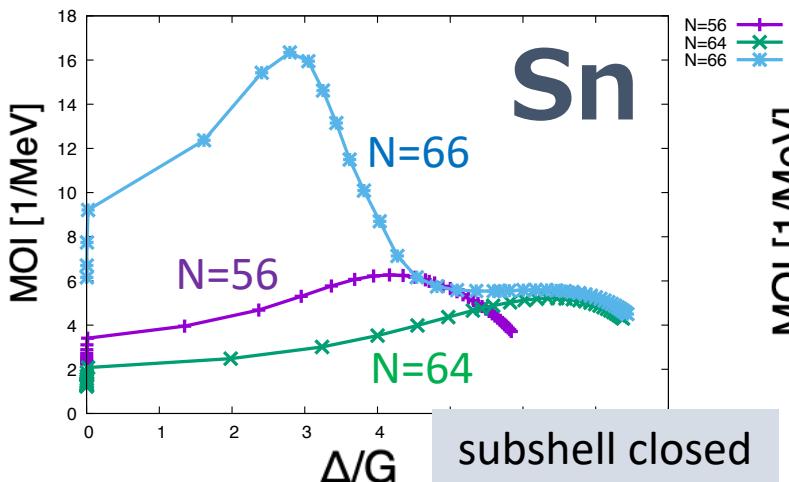
✖ When pairing strength $G \sim 0$ (NO pairing)

Ex)



Subshell closed

NOT Subshell closed



Sn: N=54~78

Ni : N=34~46

$$\Delta/G = \sum_{\nu} \Omega_{\nu} U_{\nu} V_{\nu}$$

Δ increase → MOI decrease
(except subshell closed nuclei)

Different from spatial rotation
 $MOI \propto \beta^2$

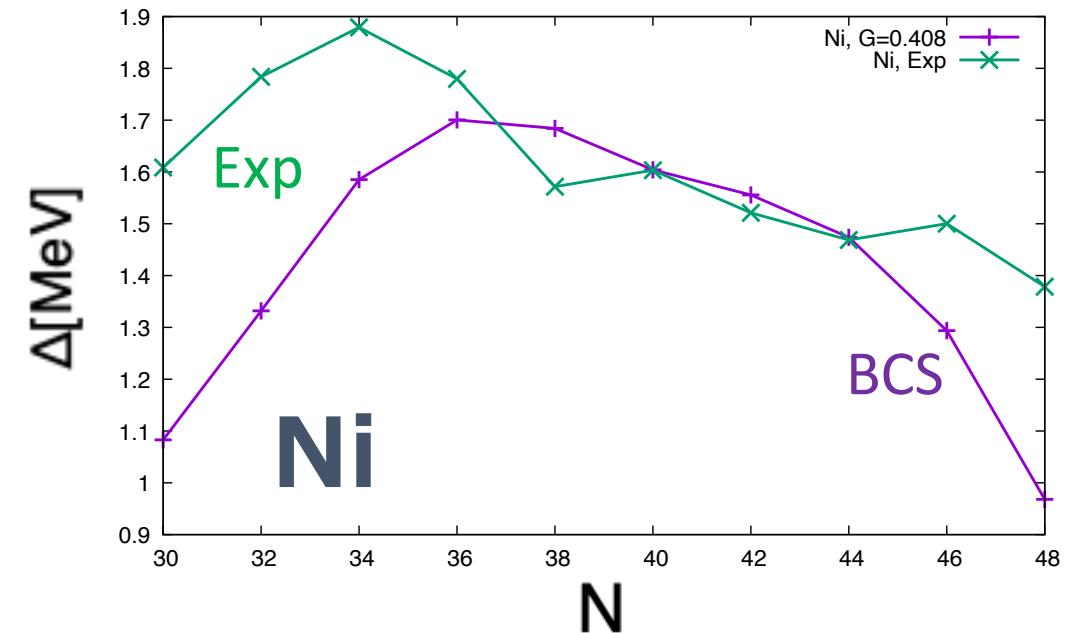
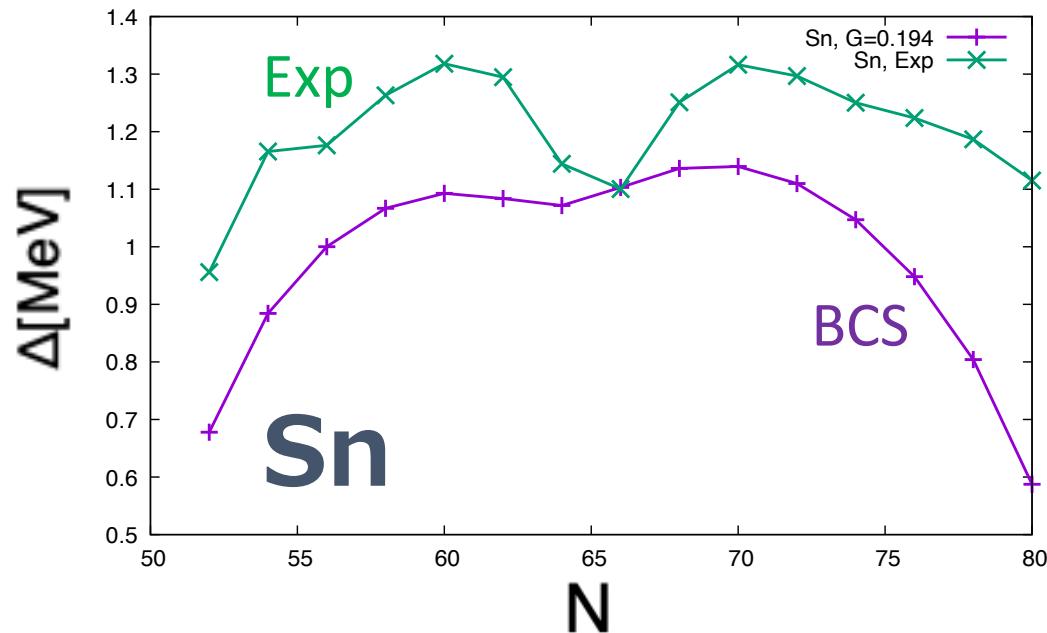
Summary & Future plans

We investigate the property of MOI of pairing rotation using BCS theory.

- In both Sn and Ni, the relationship between particle number and the MOI of pairing rotation is **roughly consistent** with that between angular momentum and MOI of **spatial rotation**.
- Particle number dependence of MOI change at ^{114}Sn and ^{68}Ni , because the orbit **most contributing to pairing collation change** into a higher energy.
- The “**deformation**” dependence of the MOI of pairing rotation is **very different** from that of **spatial rotation**.
- The “**deformation**” dependence of MOI of pairing rotation in the nuclei that are in the **subshell-closed nuclei** is different from that of in **open-shell nuclei** .

- I want to analyze the change of inner structure with respect to particle number to investigate whether there is a phenomenon that is **similar to “back bending”** in spatial rotation.
- I will check whether the same results are produced by the **HFB** and the shell model code “**KHELL**”.  now doing !

Supplement : N-Pairing gap Δ

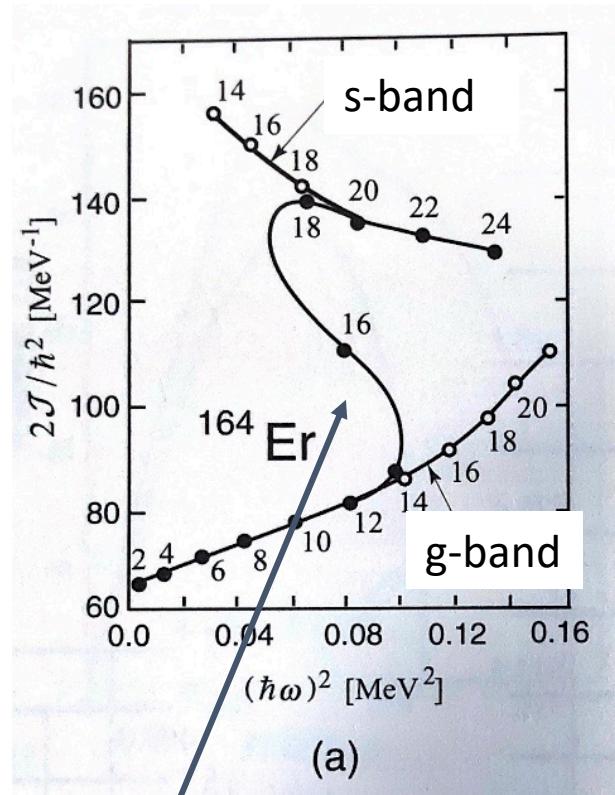


$$\text{Pairing gap} : \Delta = G \sum_{\nu>0} \Omega_\nu \Delta_\nu$$

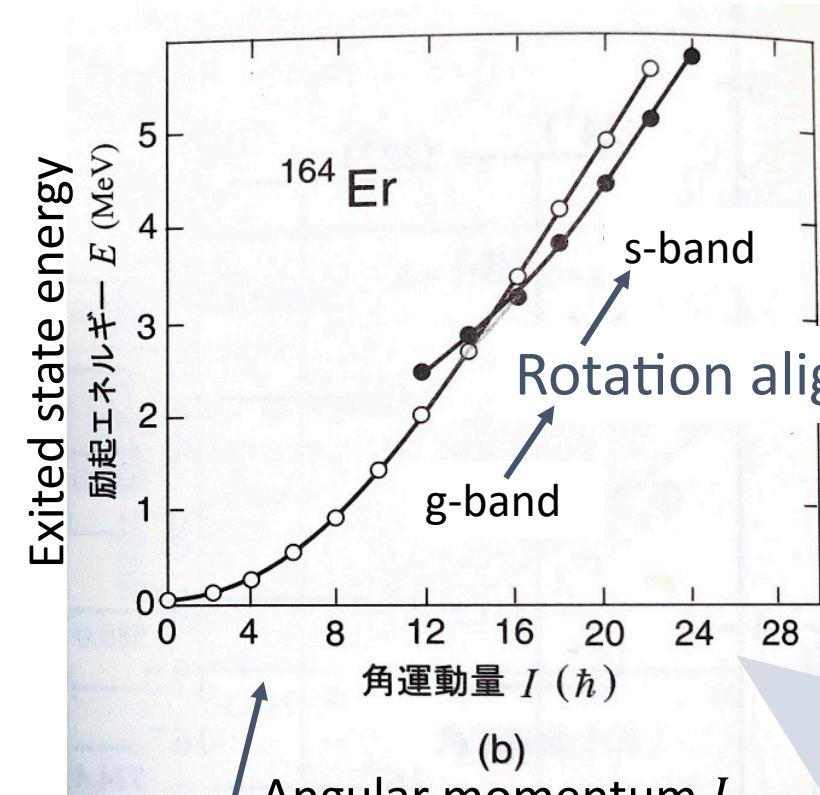
$$= G \sum_{\nu>0} \Omega_\nu \langle a_\nu a_{\bar{\nu}} \rangle$$

👉 Energy because of pairing
Something like “deformation”

Supplement : Back-bending



Back-bending



Band crossing

Back-bending :

MOI suddenly increases at angular momentum $I=14,16$ with respect to angular velocity ω

because

Rotation alignment

Band-crossing :

Rotational energy jump from g-band to s-band

Whether there is a similar phenomenon in pairing rotation? (N-MOI)

