

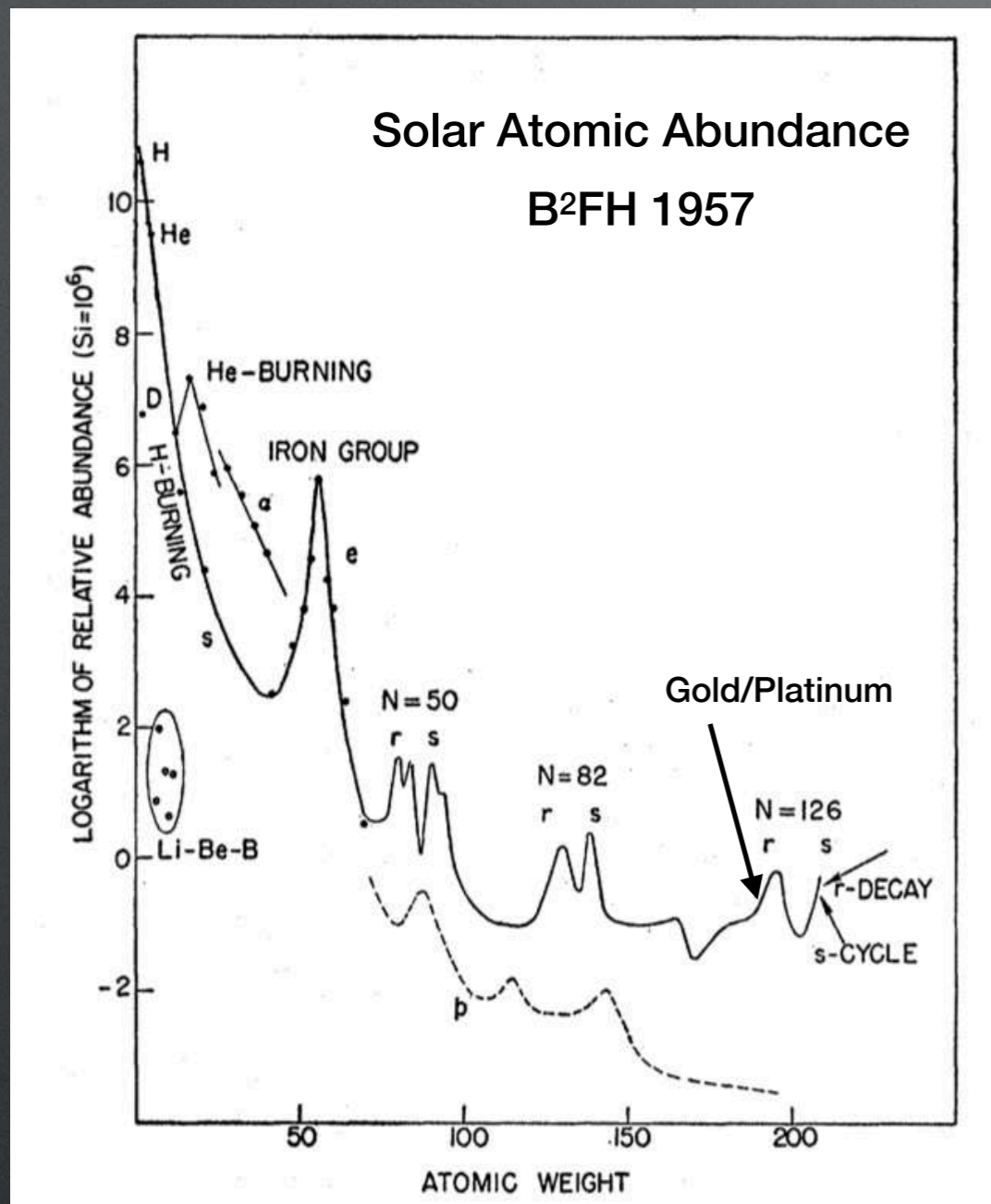
# 中性子星合体におけるR過程元素合成

仏坂健太 (東京大学)

Collaboration with M. Tanaka, D. Kato, G. Gaigalas, Y. Tarumi, N. Domoto, E. Nakar, O. Gottlieb, S. Nissanke, T. Piran, A. Horesh, A. Corsi, K. P. Mooley, A. T. Deller, M. Shibata, K. Kawaguchi, S. Fujibayashi, K. Kiuchi, P. Beniamini, R. Sari, M. M. Kasliwal, G. Hallinan



# Neutron Star Merger: heavy element factory

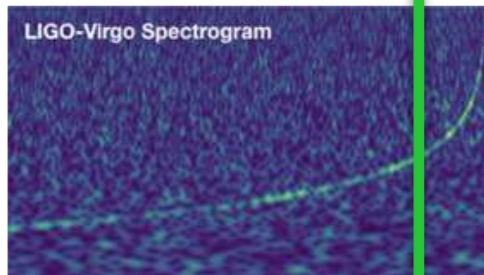


- Heavy elements (>the iron peak) are produced via neutron capture processes, (r)apid or (s)low process (B2FH 57, Cameron 57).
- The origin of r-process elements is still a mystery in astrophysics.
- The leading scenarios are core-collapse supernovae and neutron star mergers (Lattimer & Schramm 74).
- Metal poor stars & dwarf galaxies suggest that the r-process events must be rare (e.g. Ji et al 2015): r-process rate  $\ll$  supernova rate.  
(この後、垂水さん)
- A kilonova in GW170817 reveals that mergers produce some r-process elements.



# The first neutron merger in 2017 and its Electromagnetic (EM) counterparts

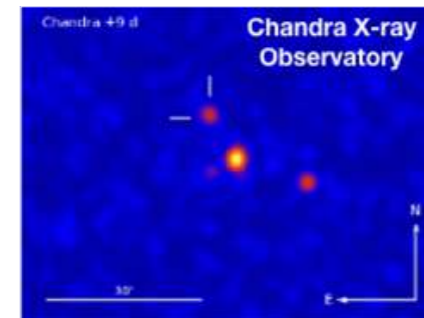
**Gravitational waves**  
(2017 Aug 17.5)  
**T = 0**



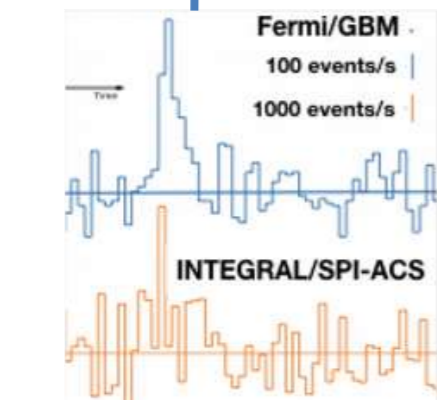
**Near Infrared**  
**T = 11h 36m**



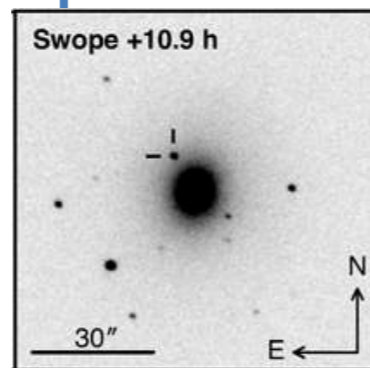
**X-rays**  
**T = 9 days**



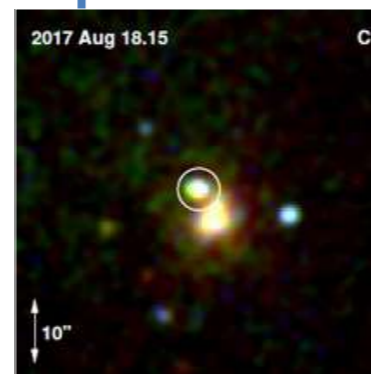
**Delayed Arrival of X-rays and Radio**



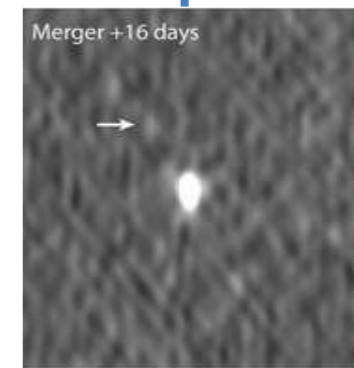
**Gamma-rays**  
**T = 1.7 seconds**



**Optical**  
**T = 10h 52m**



**Ultraviolet**  
**T = 15 hours**



**Radio**  
**T = 16 days**

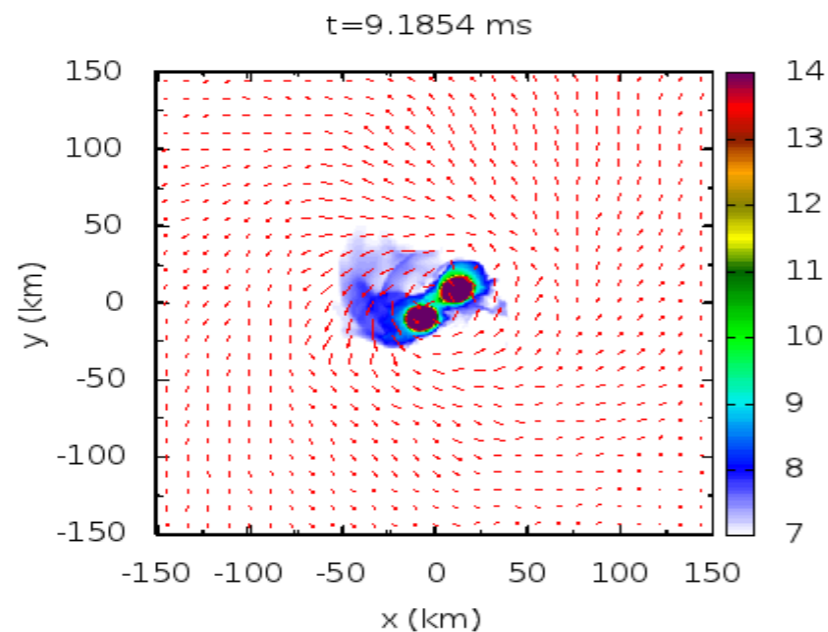
**Kilonova**

**Afterglow**

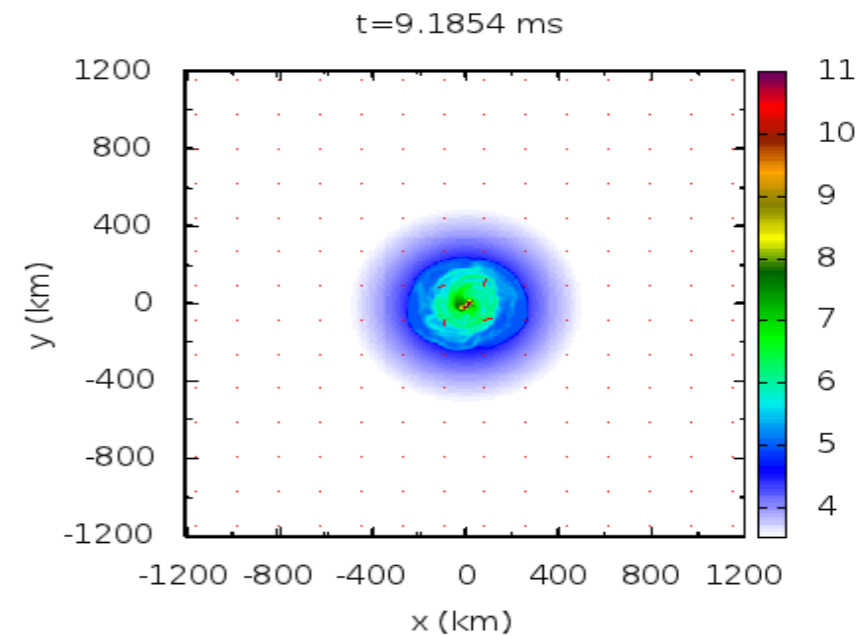
# Dynamical Ejecta in Merger

KH + 13

300 km x 300 km



2400 km x 2400 km

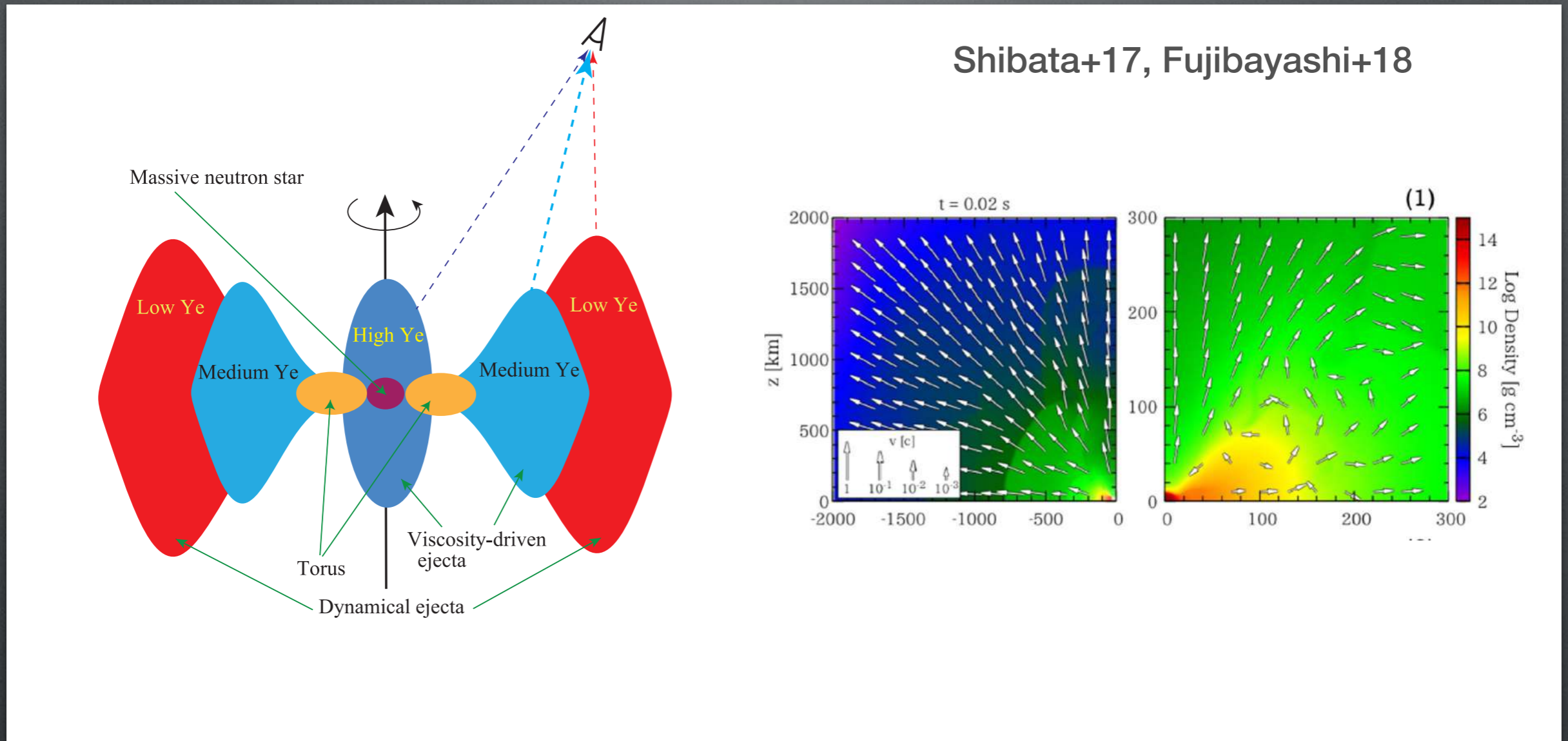


- Dynamical ejecta mass  $\sim 0.01 M_{\text{sun}}$ ,  $v \sim 0.2-0.8c$  driven by tidal force and shock heating.
- $E \sim 10^{51}$  erg, resulting in EM emission.

also Baustein + 13, Piran + 13, Rosswog 2013, Kyutoku+15, Sekiguchi + 15, 16, Radice+16



# Disk Outflow from the remnant



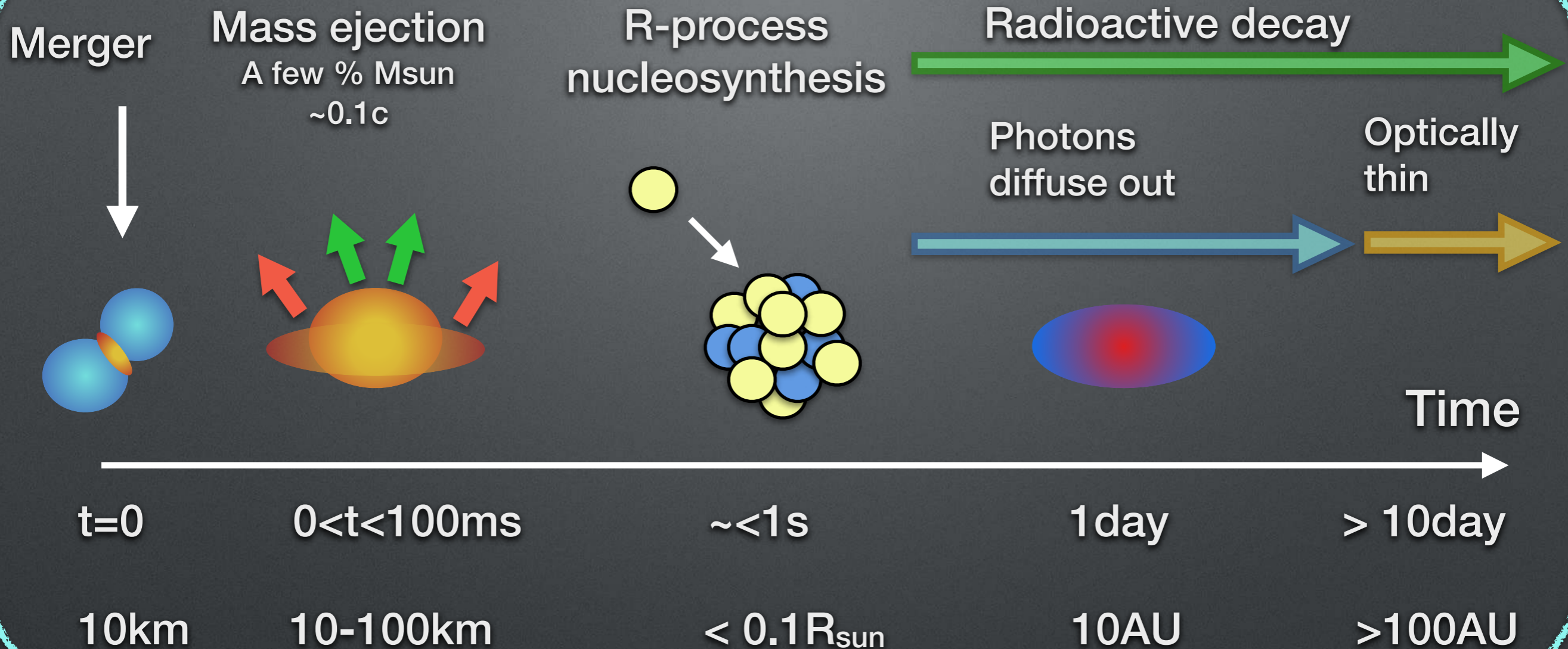
also Baustein + 13, Piran + 13, Rosswog 2013, Kyutoku+15, Sekiguchi + 15, 16, Radice+16

- Disk ejecta mass can be  $\sim 0.05M_{\text{sun}}$ , but slower  $v \sim 0.1c$ .
- In total, a few %  $M_{\text{sun}}$  and  $E \sim 10^{51}$  erg, resulting in EM emission.



# Kilonova

Li & Paczynski 98, Kulkarni 05, Metzger + 10, Barnes & Kasen 13, Tanaka & KH 13

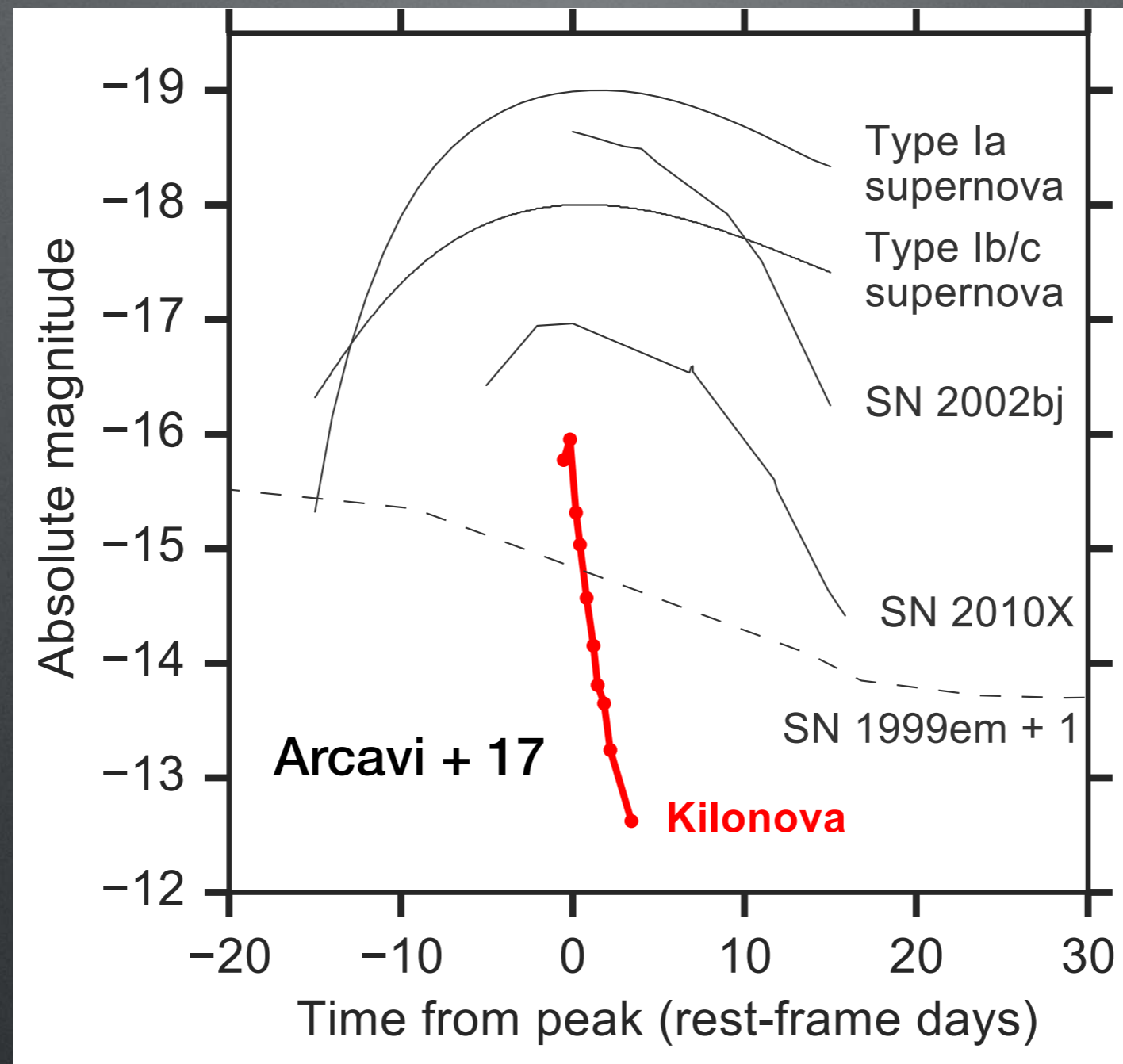


- Powered by radioactivity of r-process nuclei.
- The peak luminosity  $\sim 10^3\text{-}10^4$  x nova.
- Spectrum  $\sim$  quasi-thermal.
- Atomic lines dominate the opacity (photon absorption).



# A Kilonova in GW170817

Arcavi+17, Coulter+17, Lipunov+17, Soares-Santos+17, Tanvir+17,  
Valenti+17, Kasliwal+17, Drout+17, Evans+17, Utsumi+17

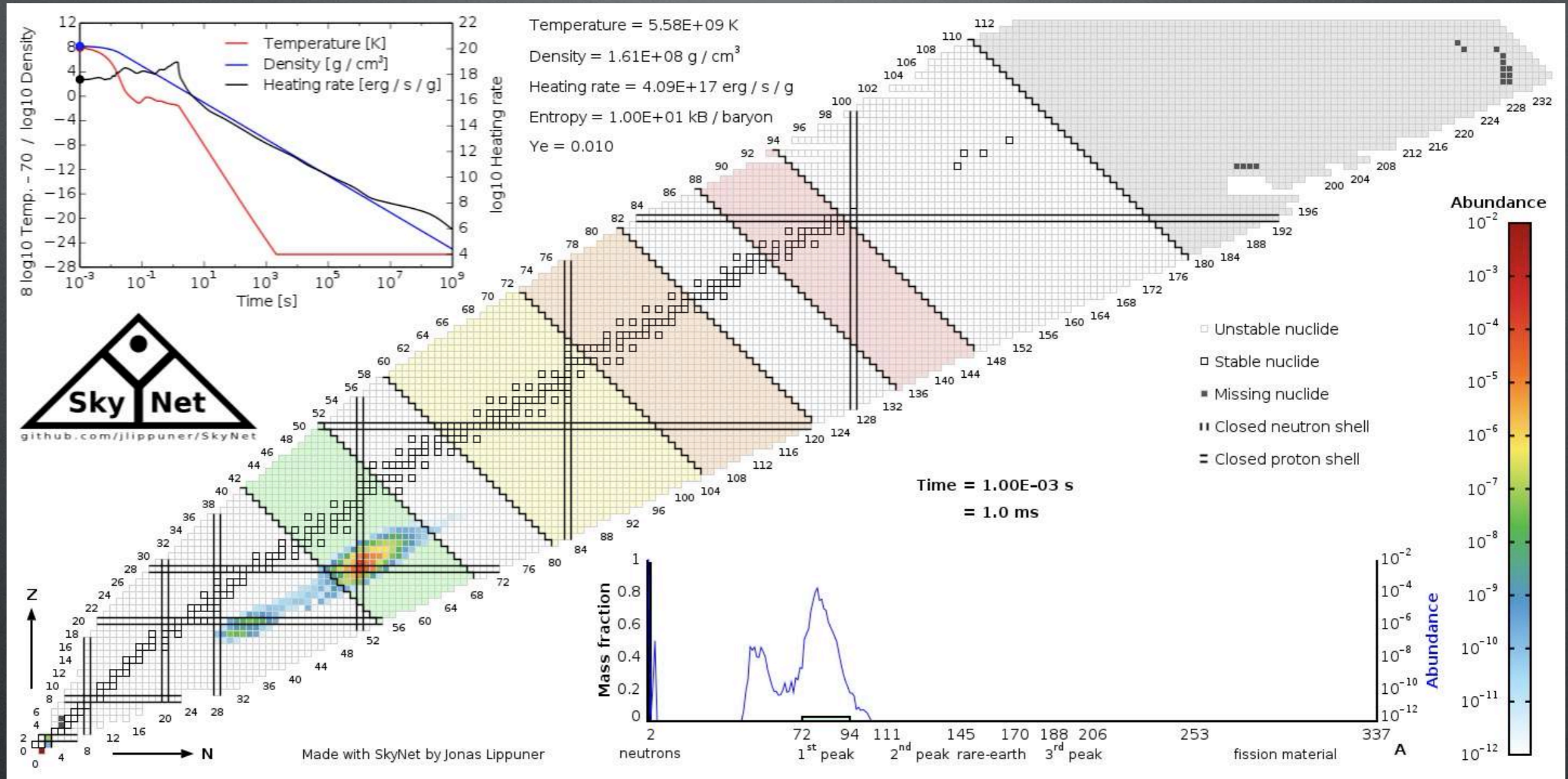


The kilonova is much fainter and faster than supernovae.

Knowing the merger times and locations from GW greatly helps to find kilonovae.



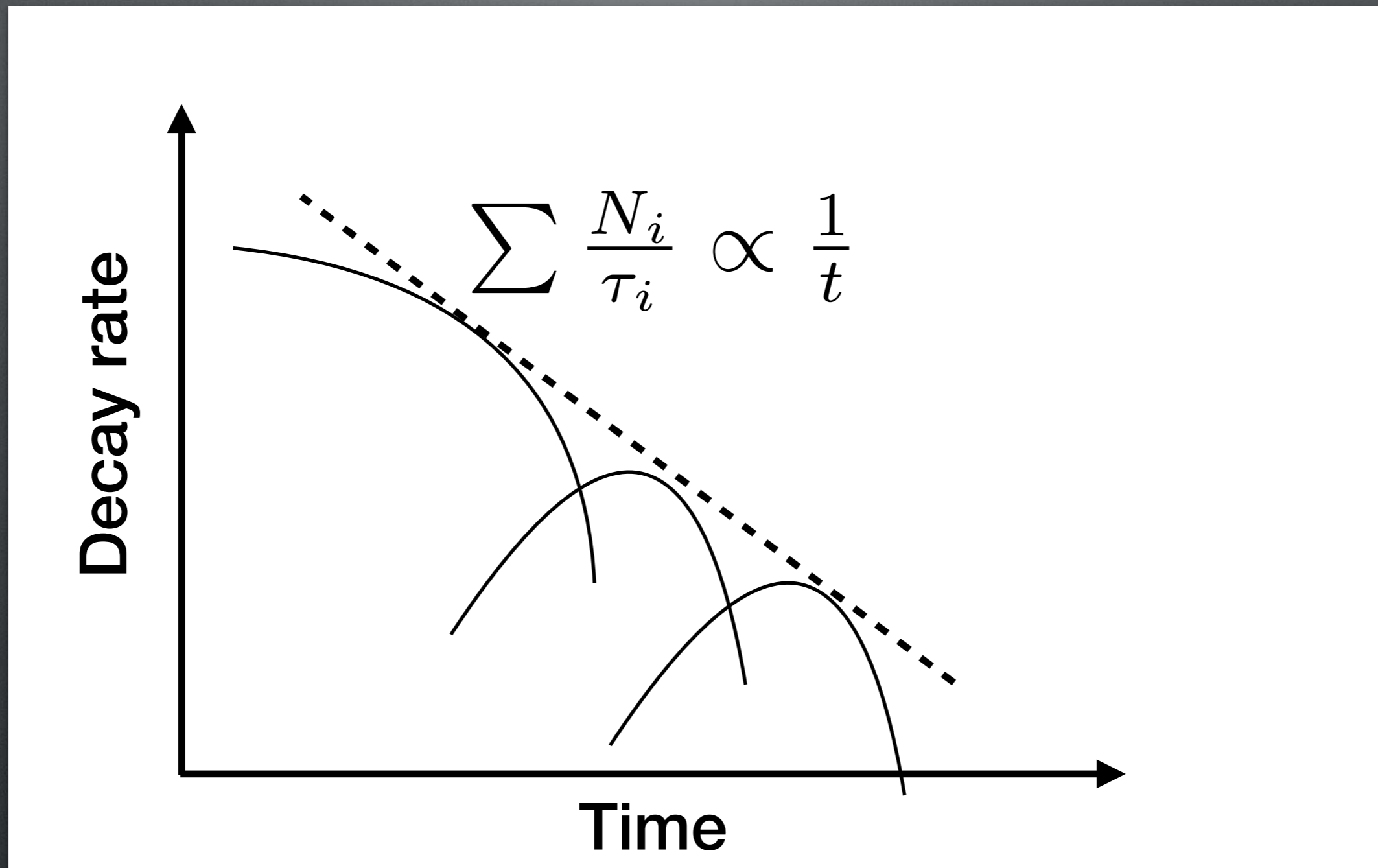
# R-process Nucleosynthesis in merger ejecta





# R-process heating

(Metzger et al 10, Goriely et al 11, Roberts et al 11, Korobkin et al 12, Wanajo et al 2014, Lippuner and Roberts 2015, KH, Sari, Piran 2017)





# Simple estimate of heating

KH,Sari,Piran 17

- Many radioactive decay chains:

$$\frac{dN}{dt} \propto \frac{1}{t} \quad \longrightarrow \quad \dot{Q}(t) \propto \frac{E(t)}{t}$$

- A relation between the lifetime and decay energy:

$$\tau \propto E^{-5} \quad \longrightarrow \quad \dot{Q}(t) \propto t^{-1.2}$$

- Physical constants in beta decay:

Energy:  $m_e c^2$

Time:  $t_F \equiv \frac{2\pi^3}{G_F^2} \frac{\hbar^7}{m_e^5 c^4} \approx 9000 \text{ s}$

Heating rate per unit mass:

$$\longrightarrow \quad \dot{Q}(t) \sim \frac{1}{\langle M \rangle} \frac{m_e c^2}{t_F} \left( \frac{t}{t_F} \right)^{-1.2} \sim 10^{10} t_{\text{day}}^{-1.2} \text{ erg/s/g}$$

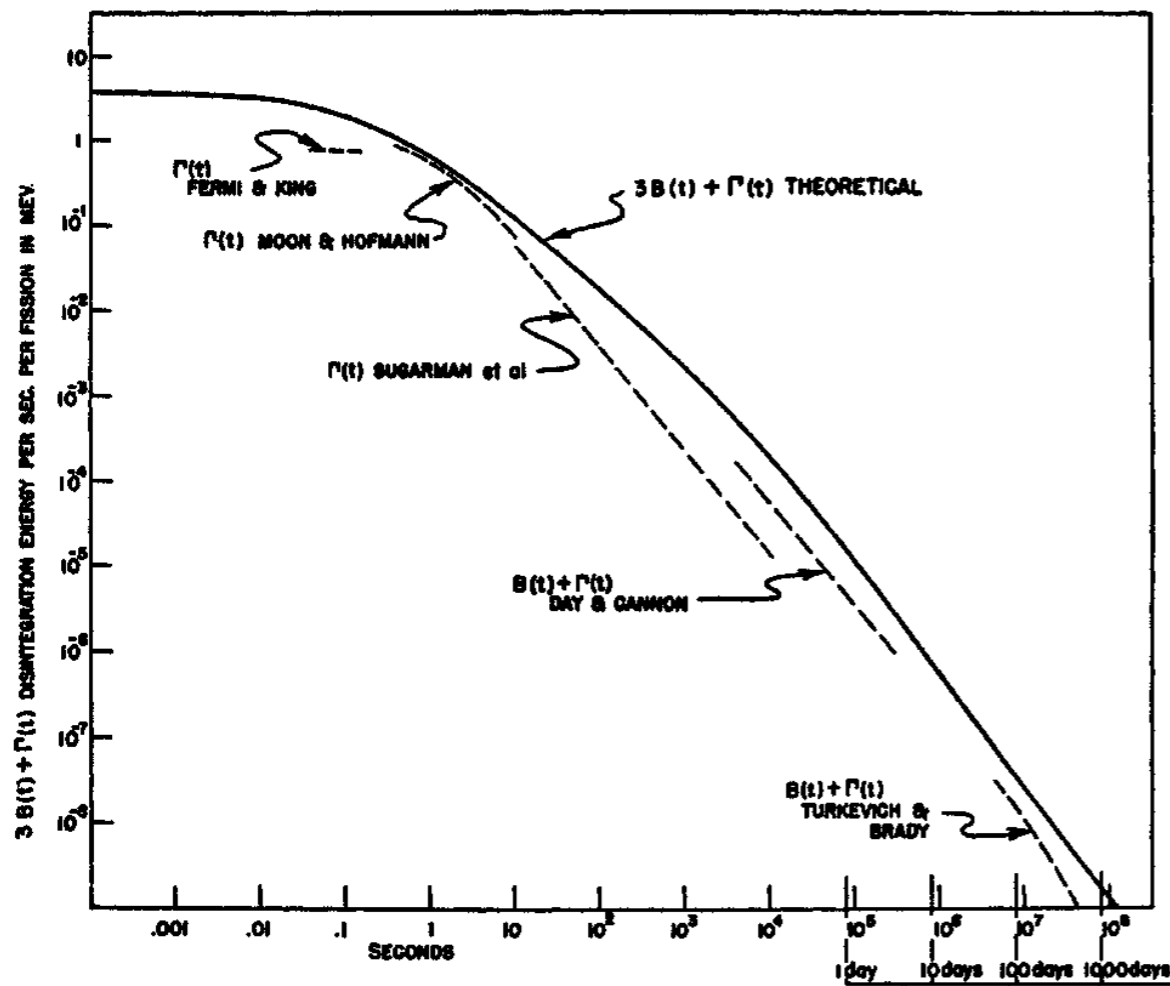
Very similar to the nuclear network results.



# Energy source: radioactive decay of many species

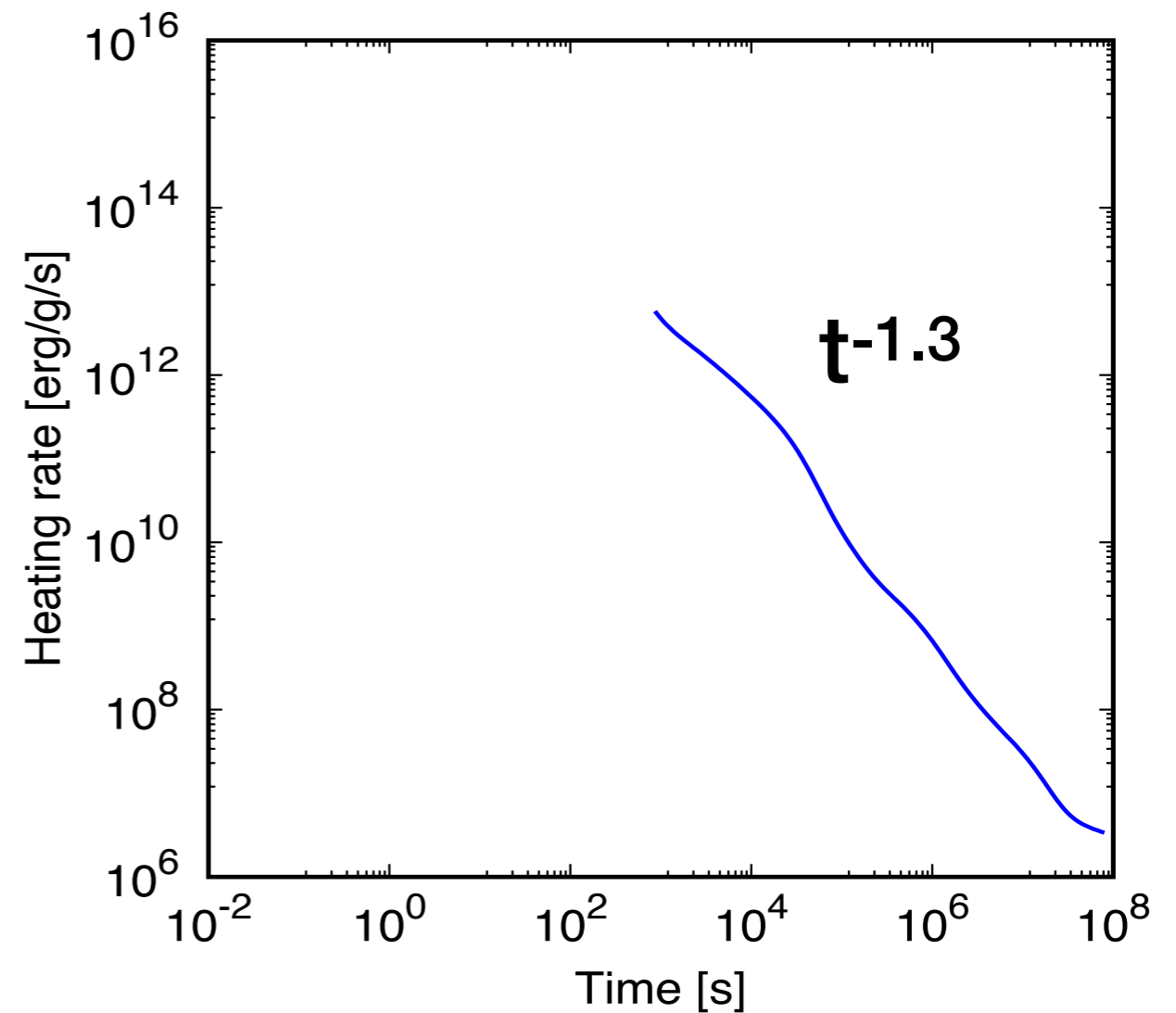
Way & Wigner 1948

KH, Sari, Piran 2017 (also Metzger + 10, Korobkin+11, Goriely+11, Roberts+11, Wanajo+14,18, Lippuner & Roberts 14)



(a)

Heating rate of nuclear waste



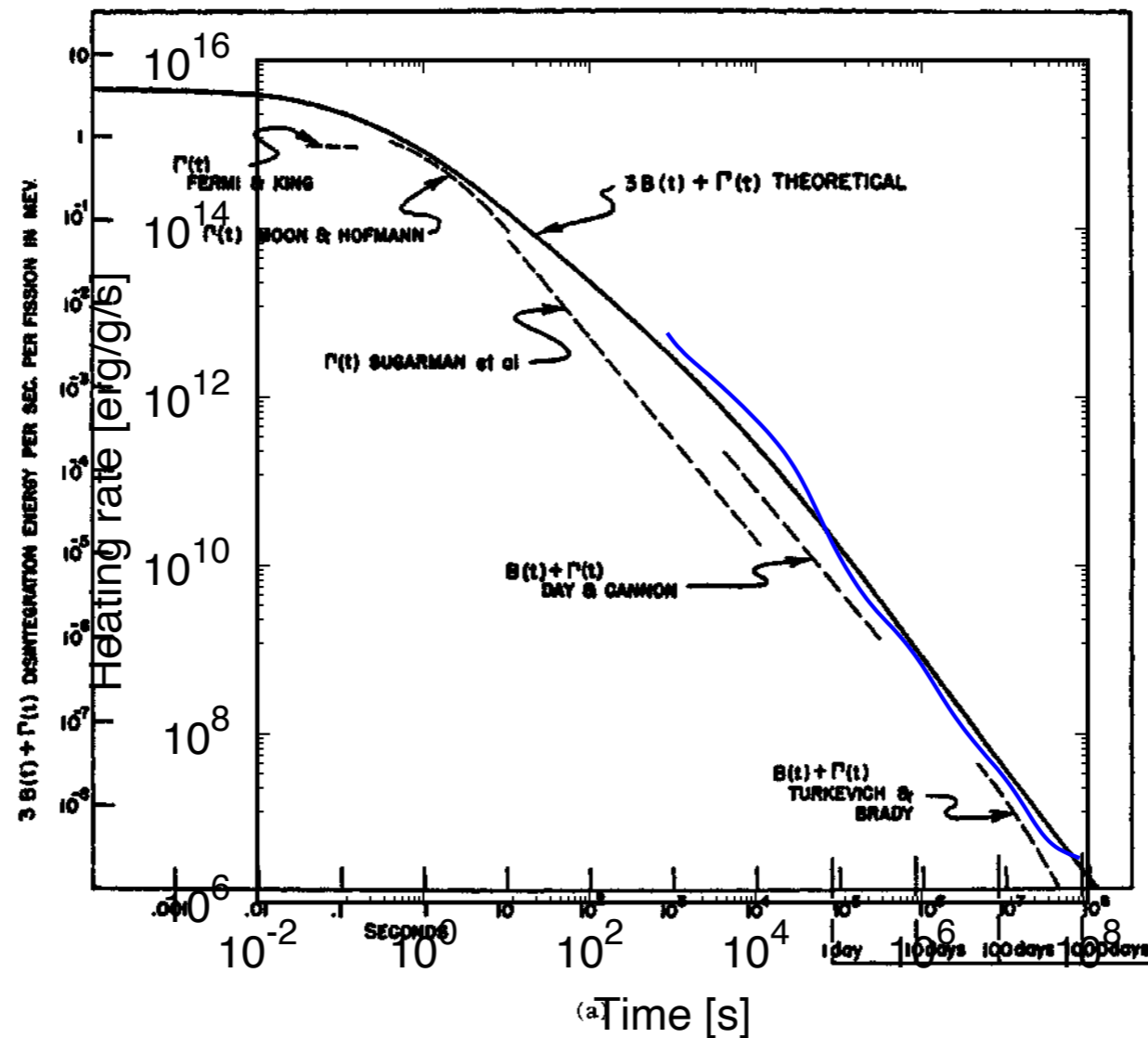
Heating rate of r-process



# Energy source: radioactive decay of many species

Way & Wigner 1948

KH, Sari, Piran 2017 (also Metzger + 10, Korobkin+11, Goriely+11, Roberts+11, Wanajo+14,18, Lippuner & Roberts 14)



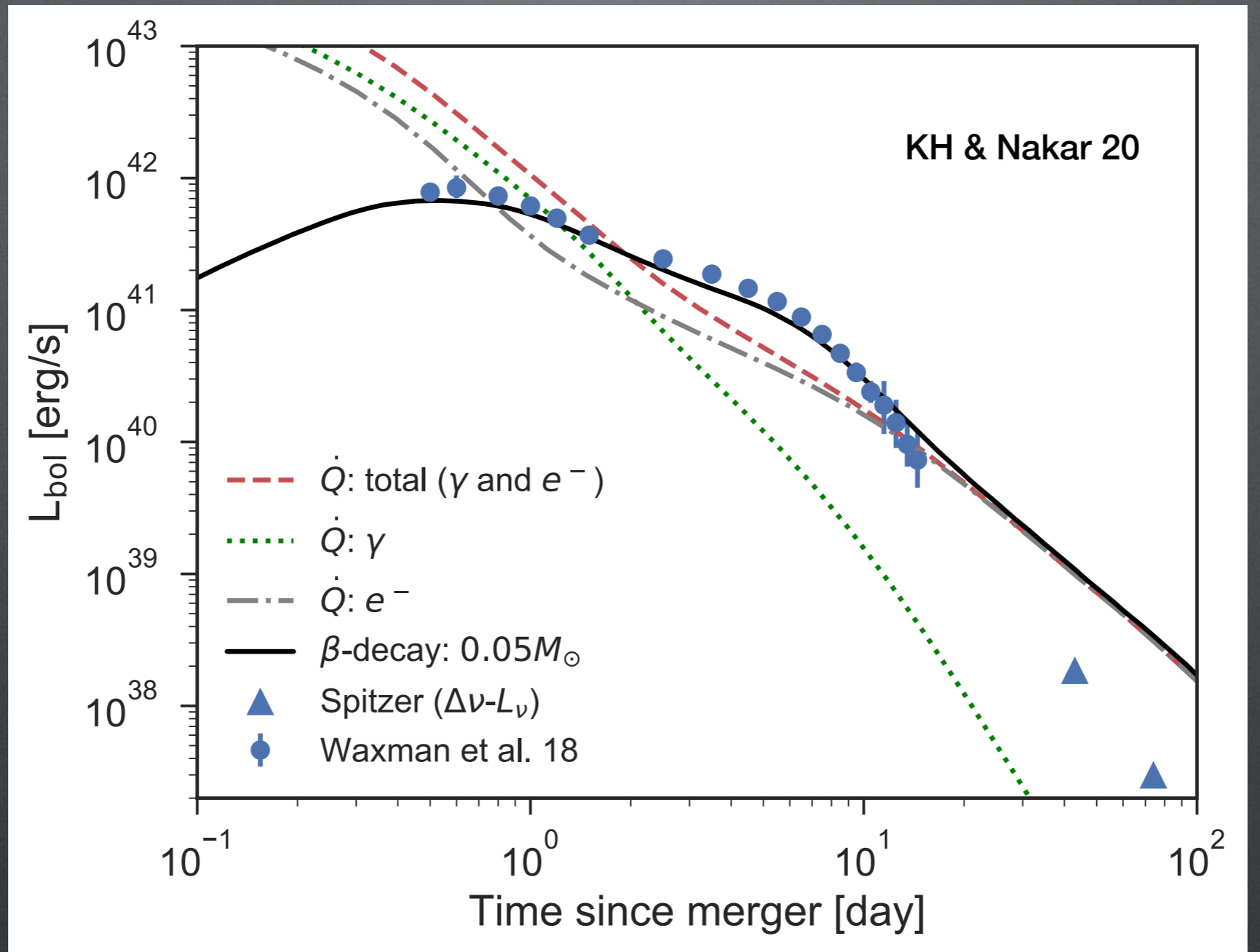
Heating rate of nuclear waste

Heating rate after r-process

This is a unique properties of the heating rates of many beta-decay chains.



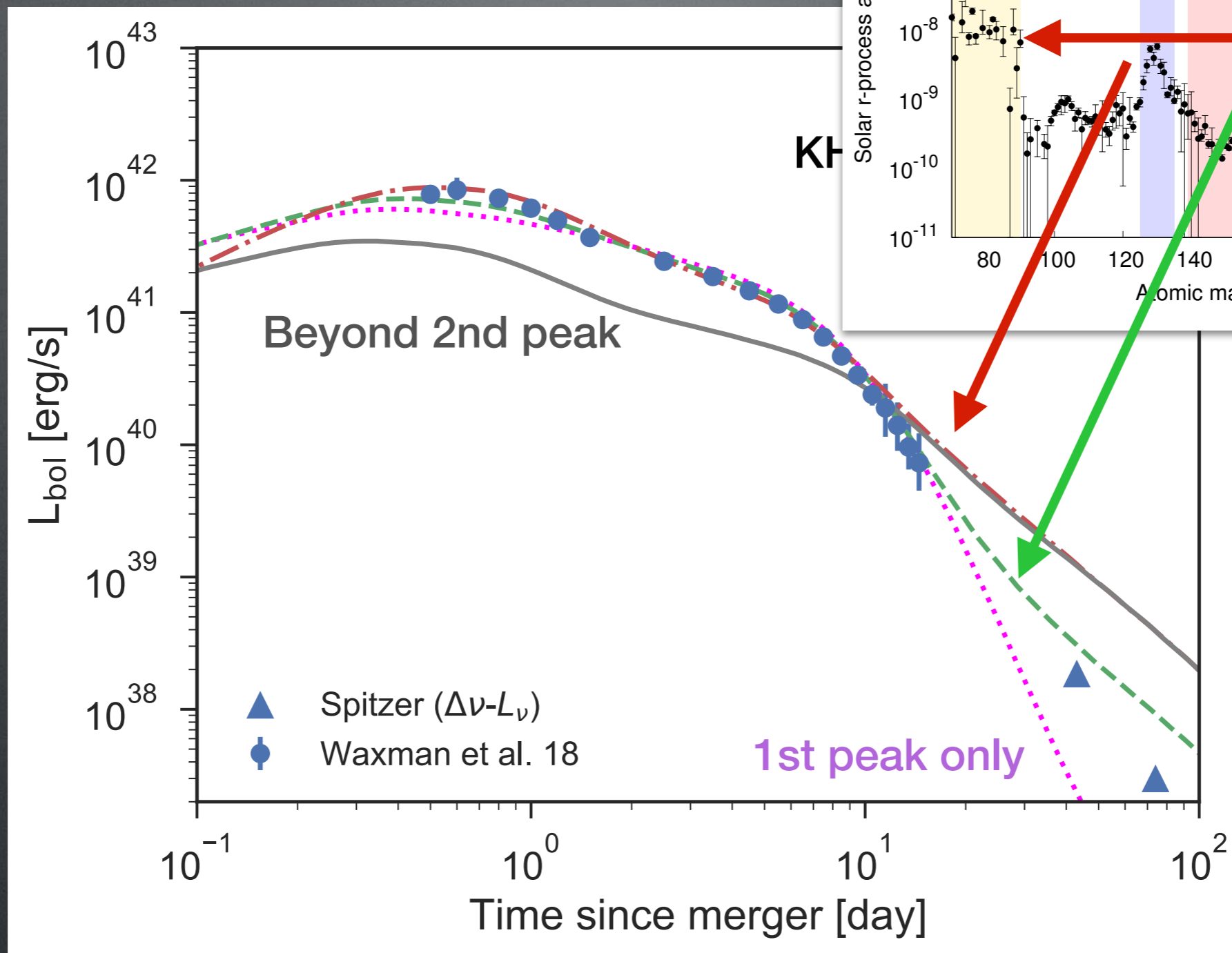
# The energy budget of the Kilonova in GW170817



- The light curve follows the  $\beta$ -decay heating rate.
- Ejecta mass is  $\sim 0.05M_{\text{sun}}$ .
- The photospheric velocity  $\sim 0.1-0.3c$ .
- The photospheric temperature evolves  $T=5000\text{K} \rightarrow 2000\text{K}$ .



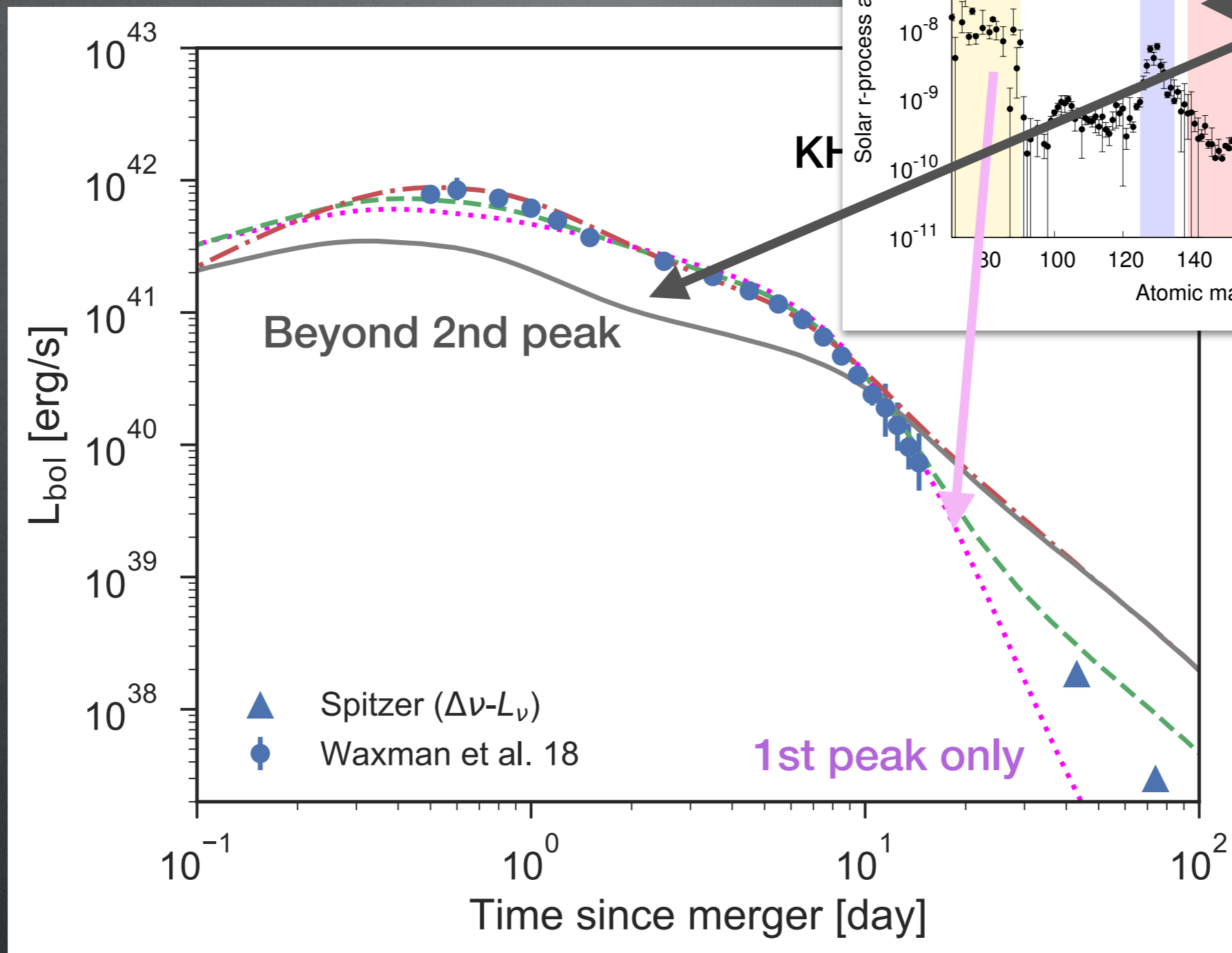
# Implication to the r-process



As long as, the 2nd and 3rd elements are abundant, radioactivity can explain the energy budget.



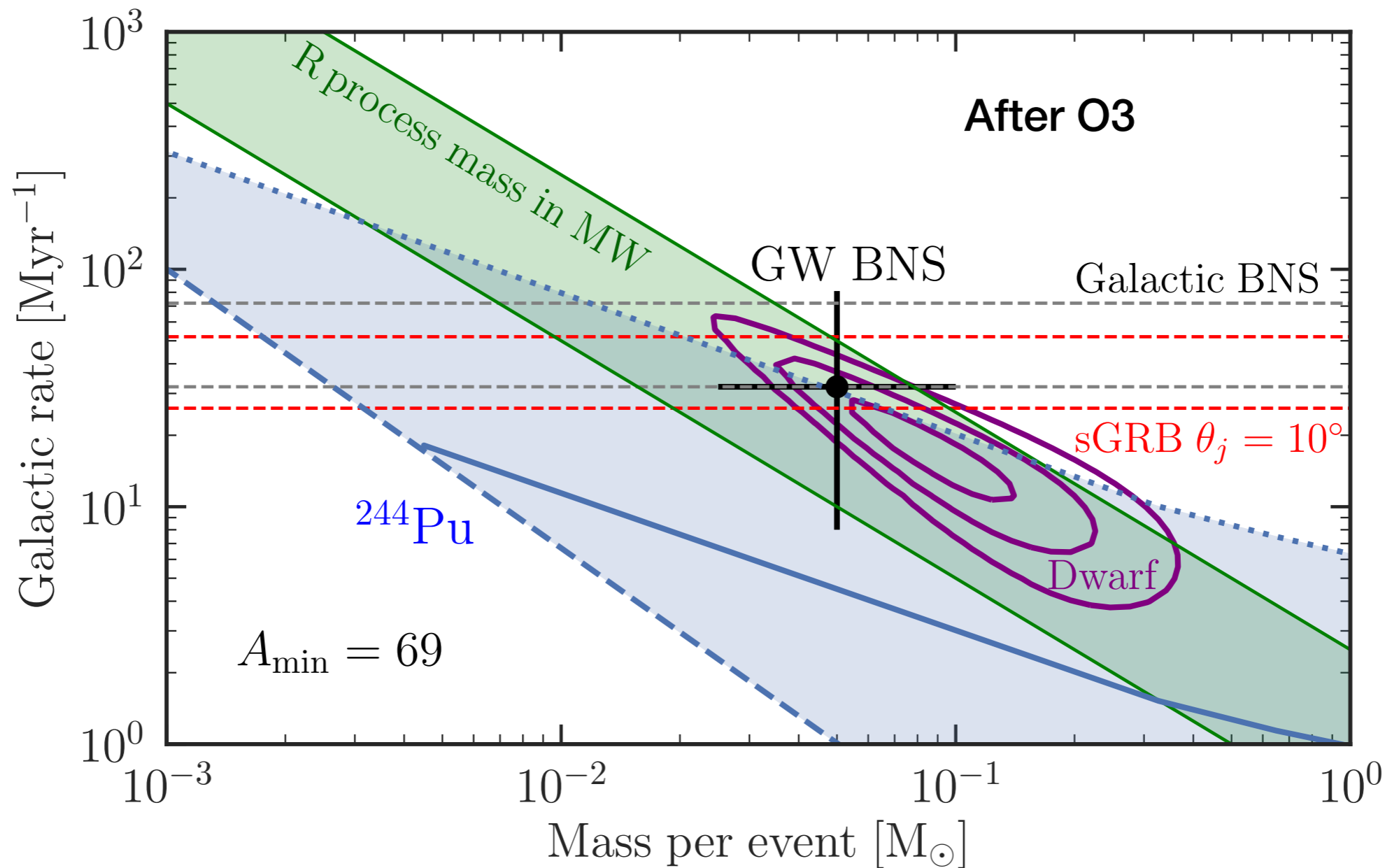
# Implication to the r-process



- (i) the 1st peak only cannot explain the observation.
- (ii) Beyond 2nd peak only also cannot.

# R-process mass budget from GWTC-2

KH, Piran, Paul 15, KH, Beniamini, Piran 18



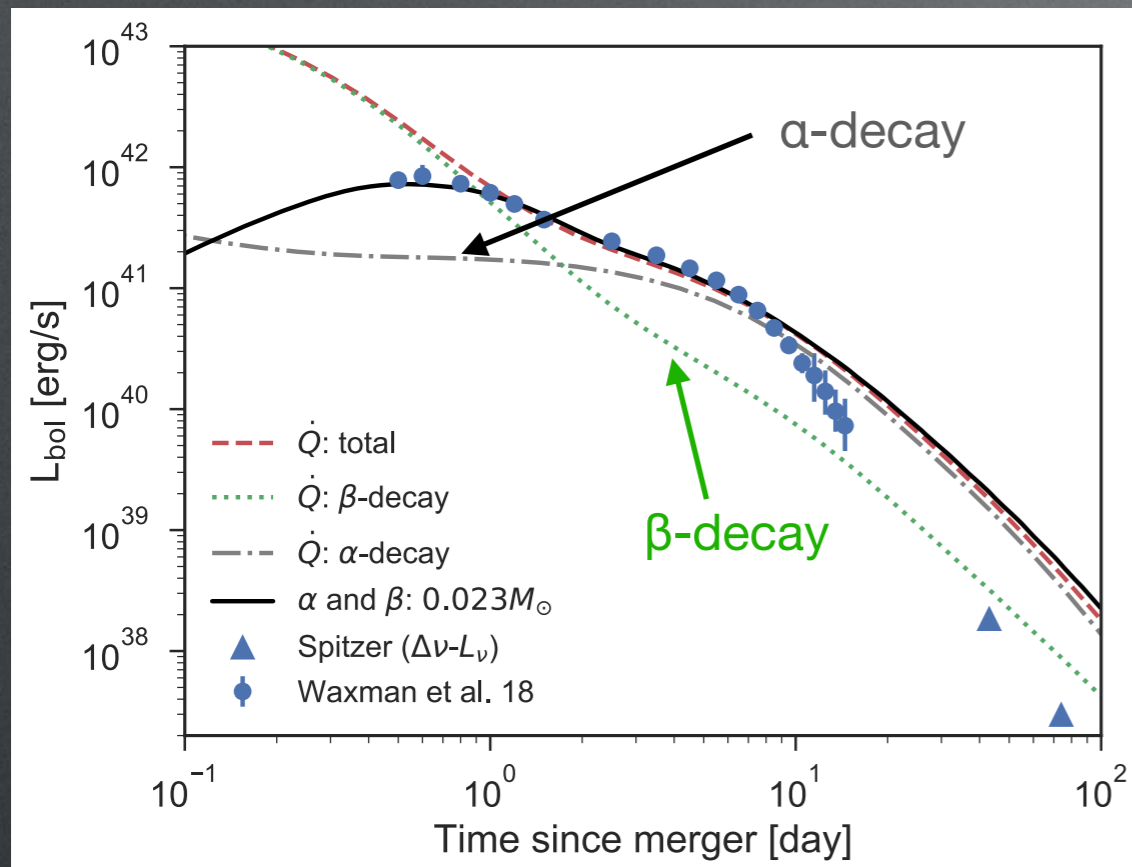
Ref: Goriely 1999, Lodders et al 2009, Wanderman & Piran 2015, Fong+2015, KH, Piran, Paul 2015, Beniamini, KH, Piran 2016, Pol, McLaughlin, Lorimer 2019, KH & Nakar 2020, LVC 2020



# $\alpha$ -decay and fission as the energy source?

KH & Nakar 20, see also, Wanajo+14, KH+16, Barnes+16, Zhu+18, Wu+19

## $\alpha$ -decay



<sup>216</sup> U	<sup>217</sup> U	<sup>218</sup> U	<sup>219</sup> U		<sup>221</sup> U	<sup>222</sup> U	<sup>223</sup> U	<sup>224</sup> U	<sup>225</sup> U	<sup>226</sup> U	<sup>227</sup> U	<sup>228</sup> U	J $\pi$	0+
													T <sub>1/2</sub> or $\Gamma$	3.6316 d 0.0023
<sup>215</sup> Pa	<sup>216</sup> Pa	<sup>217</sup> Pa	<sup>218</sup> Pa	<sup>219</sup> Pa	<sup>220</sup> Pa	<sup>221</sup> Pa	<sup>222</sup> Pa	<sup>223</sup> Pa	<sup>224</sup> Pa	<sup>225</sup> Pa	<sup>226</sup> Pa	<sup>227</sup> Pa	Delta (keV)	18825.832 1811
													Bind/A (keV)	7679.9236 81
<sup>214</sup> Th	<sup>215</sup> Th	<sup>216</sup> Th	<sup>217</sup> Th	<sup>218</sup> Th	<sup>219</sup> Th	<sup>220</sup> Th	<sup>221</sup> Th	<sup>222</sup> Th	<sup>223</sup> Th	<sup>224</sup> Th	<sup>225</sup> Th	<sup>226</sup> Th	Mass ( $\mu$ AMU)	224020210.361 1944
													Q $\alpha$ (keV)	5788.9231 1526
<sup>213</sup> Ac	<sup>214</sup> Ac	<sup>215</sup> Ac	<sup>216</sup> Ac	<sup>217</sup> Ac	<sup>218</sup> Ac	<sup>219</sup> Ac	<sup>220</sup> Ac	<sup>221</sup> Ac	<sup>222</sup> Ac	<sup>223</sup> Ac	<sup>224</sup> Ac	<sup>225</sup> Ac	Q $\beta$ (keV)	-1408.3152 40869
													Q $\epsilon$ (keV)	-2922.782 11
<sup>212</sup> Ra	<sup>213</sup> Ra	<sup>214</sup> Ra	<sup>215</sup> Ra	<sup>216</sup> Ra	<sup>217</sup> Ra	<sup>218</sup> Ra	<sup>219</sup> Ra	<sup>220</sup> Ra	<sup>221</sup> Ra	<sup>222</sup> Ra	<sup>223</sup> Ra	<sup>224</sup> Ra	Sn (keV)	6478.7305 22609
													Sp (keV)	6845.4679 21130
<sup>211</sup> Fr	<sup>212</sup> Fr	<sup>213</sup> Fr	<sup>214</sup> Fr	<sup>215</sup> Fr	<sup>216</sup> Fr	<sup>217</sup> Fr	<sup>218</sup> Fr	<sup>219</sup> Fr	<sup>220</sup> Fr	<sup>221</sup> Fr	<sup>222</sup> Fr	<sup>223</sup> Fr	Decay	$\alpha$ 100%
														<sup>12</sup> C 4.0E-9%
<sup>210</sup> Rn	<sup>211</sup> Rn	<sup>212</sup> Rn	<sup>213</sup> Rn	<sup>214</sup> Rn	<sup>215</sup> Rn	<sup>216</sup> Rn	<sup>217</sup> Rn	<sup>218</sup> Rn	<sup>219</sup> Rn	<sup>220</sup> Rn	<sup>221</sup> Rn	<sup>222</sup> Rn	Major radiations	
													Type keV	%
<sup>209</sup> At	<sup>210</sup> At	<sup>211</sup> At	<sup>212</sup> At	<sup>213</sup> At	<sup>214</sup> At	<sup>215</sup> At	<sup>216</sup> At	<sup>217</sup> At	<sup>218</sup> At	<sup>219</sup> At	<sup>220</sup> At	<sup>221</sup> At	$\alpha$	5685.37 94.92
													$\gamma$	5448.6 5.06
<sup>208</sup> Po	<sup>209</sup> Po	<sup>210</sup> Po	<sup>211</sup> Po	<sup>212</sup> Po	<sup>213</sup> Po	<sup>214</sup> Po	<sup>215</sup> Po	<sup>216</sup> Po	<sup>217</sup> Po	<sup>218</sup> Po	<sup>219</sup> Po	<sup>220</sup> Po		240.986 4.10
														10.132 -18.044 0.371
<sup>207</sup> Bi	<sup>208</sup> Bi	<sup>209</sup> Bi	<sup>210</sup> Bi	<sup>211</sup> Bi	<sup>212</sup> Bi	<sup>213</sup> Bi	<sup>214</sup> Bi	<sup>215</sup> Bi	<sup>216</sup> Bi	<sup>217</sup> Bi	<sup>218</sup> Bi	<sup>219</sup> Bi		
<sup>206</sup> Pb	<sup>207</sup> Pb	<sup>208</sup> Pb	<sup>209</sup> Pb	<sup>210</sup> Pb	<sup>211</sup> Pb	<sup>212</sup> Pb	<sup>213</sup> Pb	<sup>214</sup> Pb	<sup>215</sup> Pb	<sup>216</sup> Pb	<sup>217</sup> Pb	<sup>218</sup> Pb	<sup>219</sup> Pb	<sup>220</sup> Pb

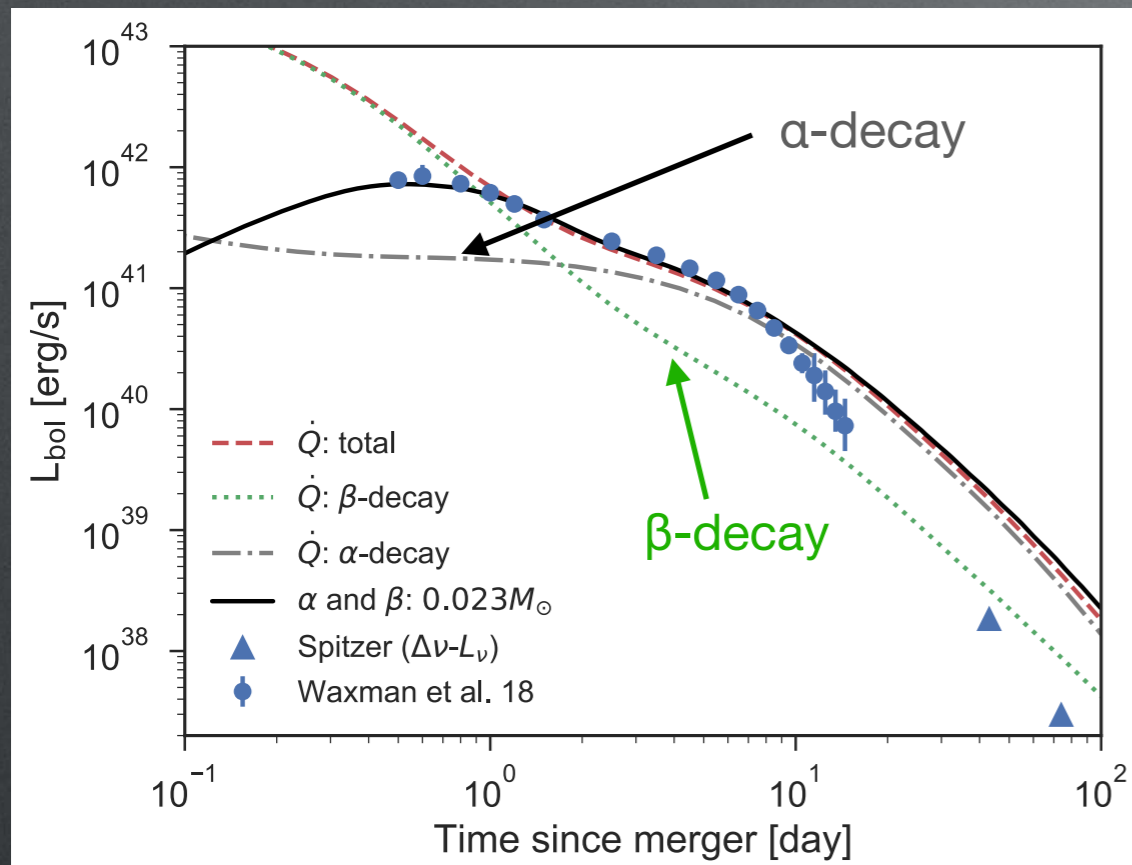
- $\alpha$ -decay can dominate the energy source (<sup>222</sup>Rn, <sup>223</sup>Ra, <sup>224</sup>Ra, <sup>225</sup>Ra). But the relative importance to  $\beta$ -decay is very uncertain.
- Spontaneous fission can also be important at late times > 10 day. But it is also highly uncertain.



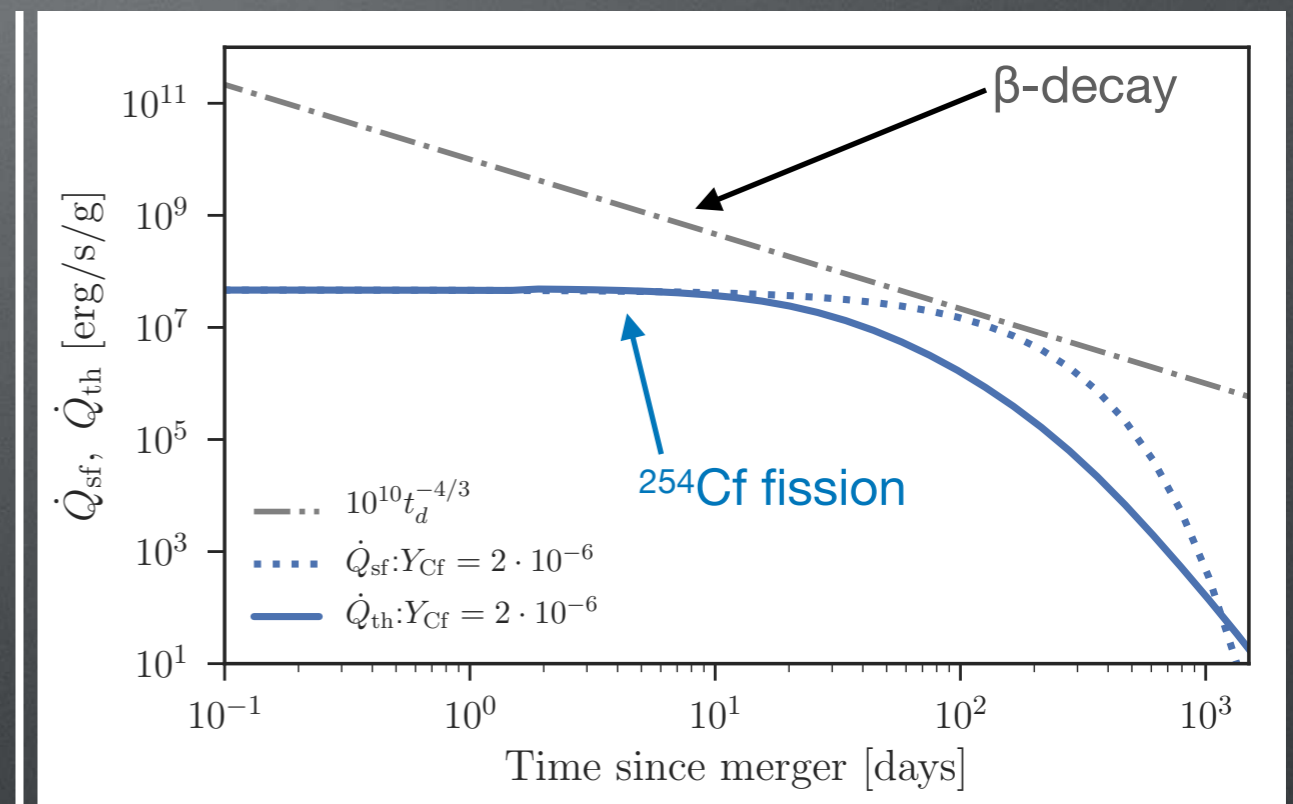
# $\alpha$ -decay and fission as the energy source?

KH & Nakar 20, see also, Wanajo+14, KH+16, Barnes+16, Zhu+18, Wu+19

$\alpha$ -decay



Spontaneous fission

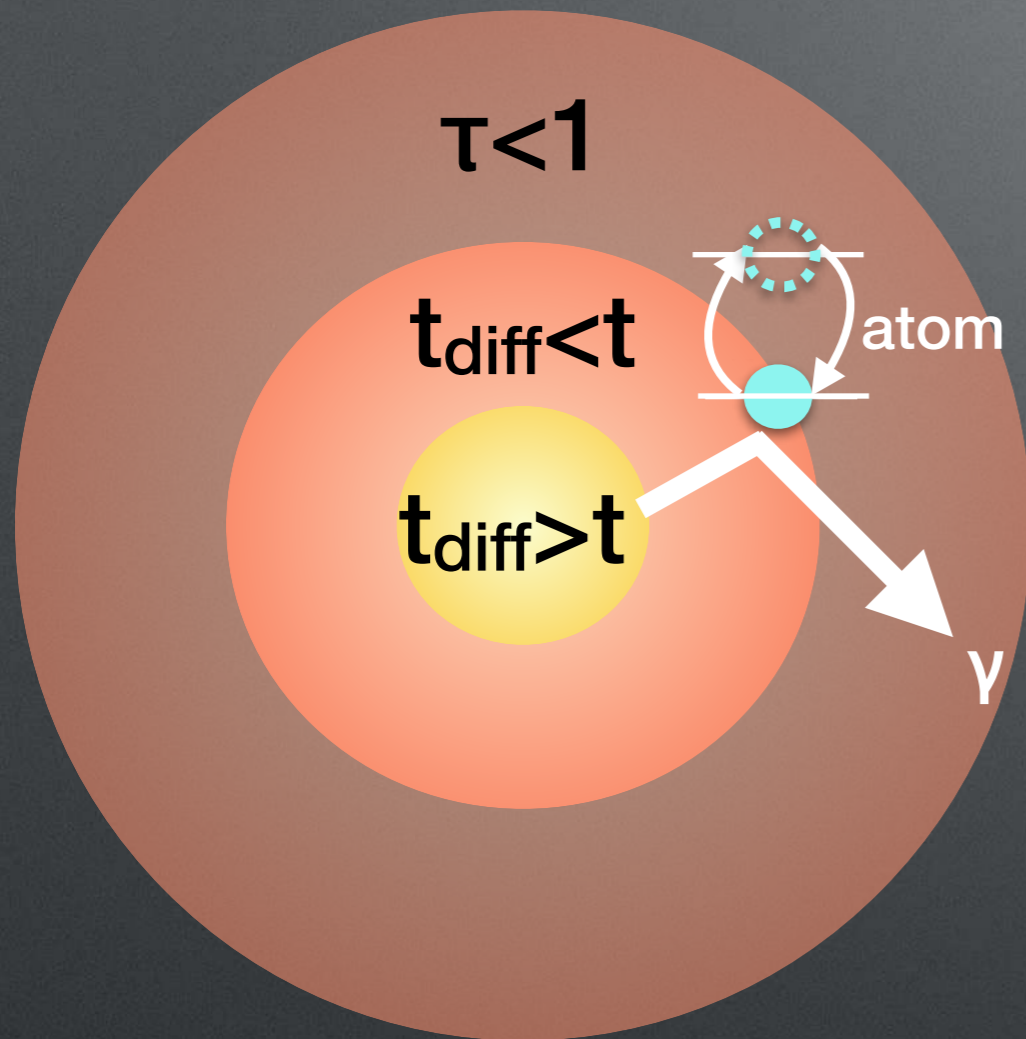


- $\alpha$ -decay can dominate the energy source ( $^{222}\text{Rn}$ ,  $^{223}\text{Ra}$ ,  $^{224}\text{Ra}$ ,  $^{225}\text{Ra}$ ). But the relative importance to  $\beta$ -decay is very uncertain.
- Spontaneous fission can also be important at late times  $> 10$  day. But it is also highly uncertain.

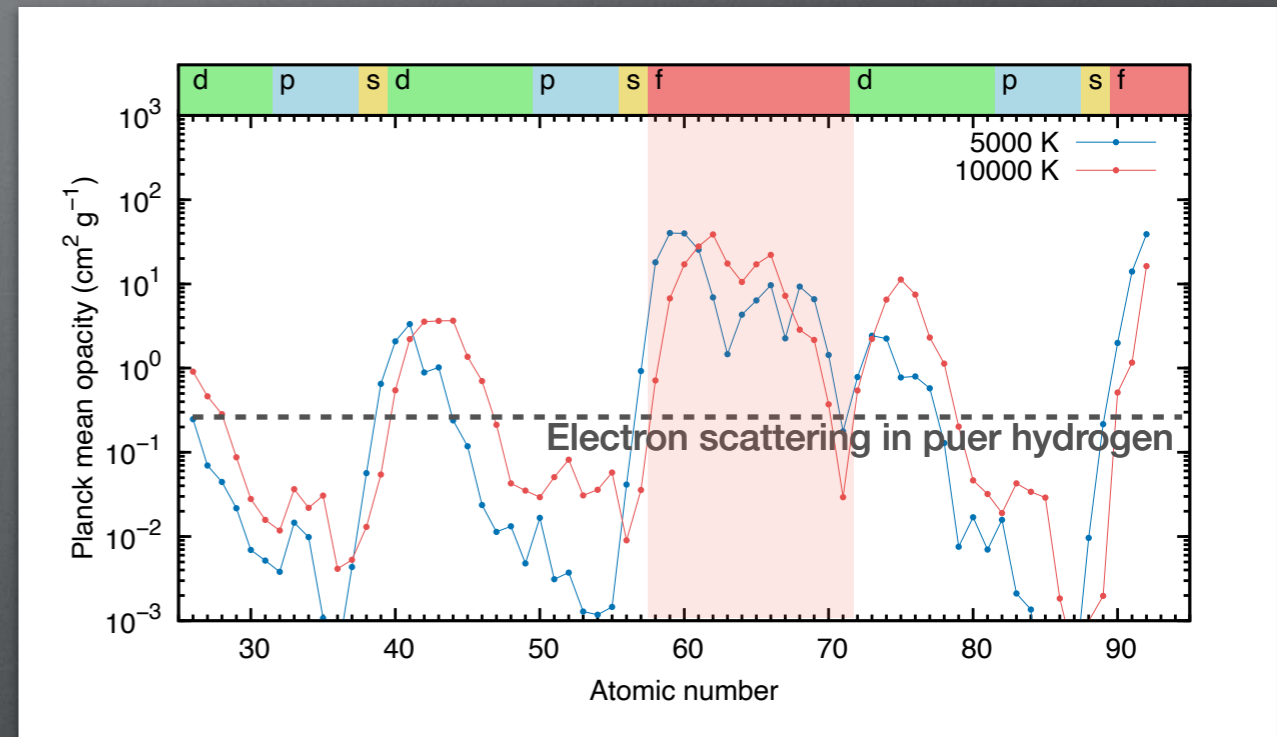


# Kilonova Radiation Transfer

Photospheric phase ( $t < 10$  day)



Tanaka+2020

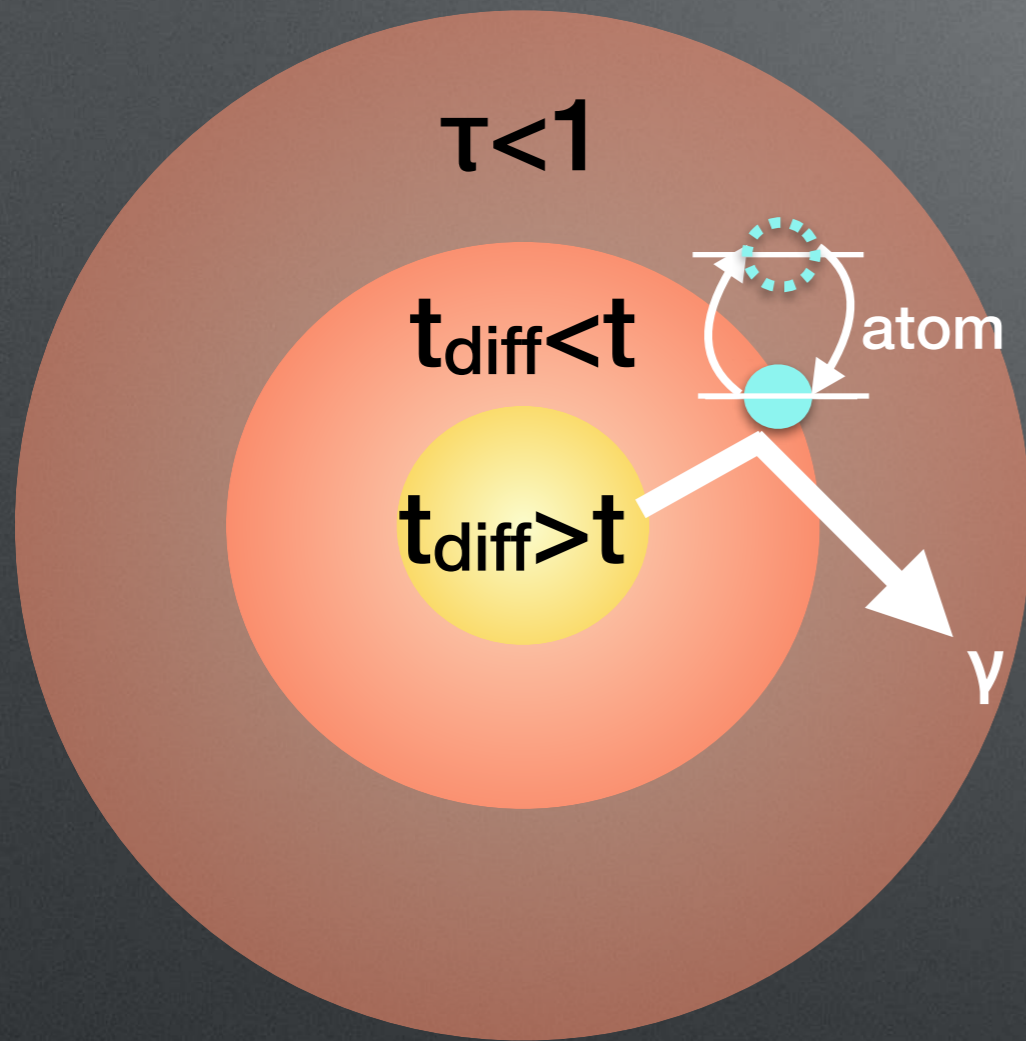


- The photosphere is in the middle ( $v \sim 0.1 - 0.2c$ ).
- The photospheric temperature is  $T \sim 10^3 \text{K}$ .
- The spectrum peaks at optical-nIR. Rich observation data exist for GW170817.

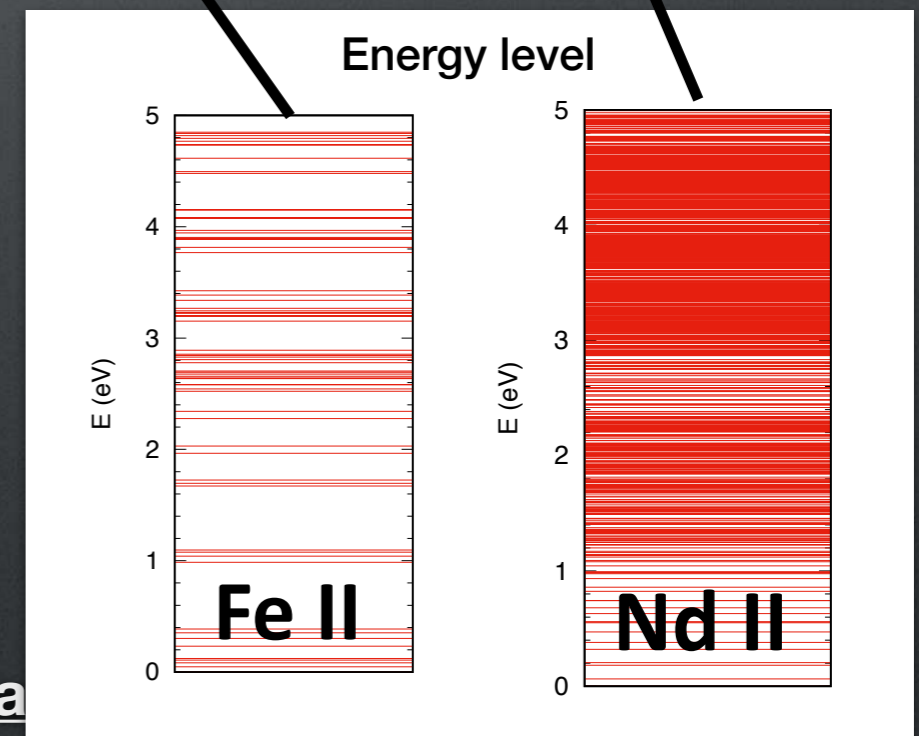
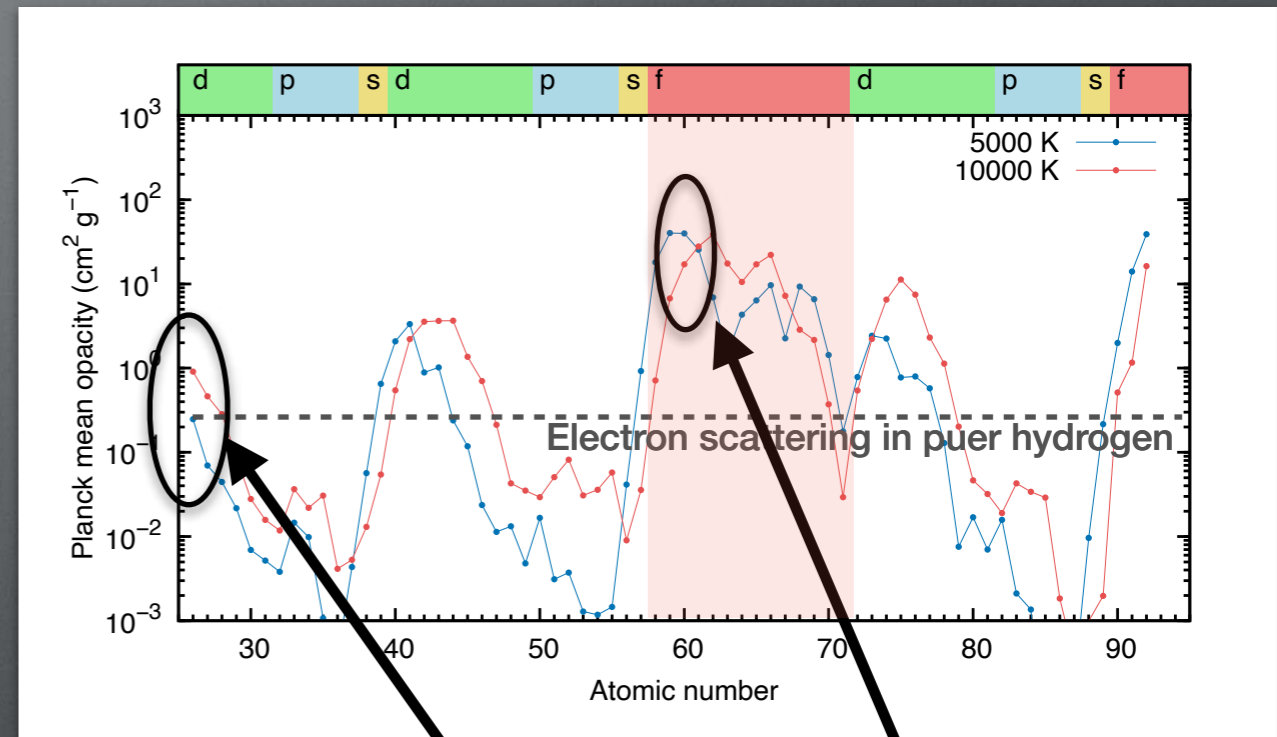


# Kilonova Radiation Transfer

Photospheric phase ( $t < 10$  day)



Tanaka+2020

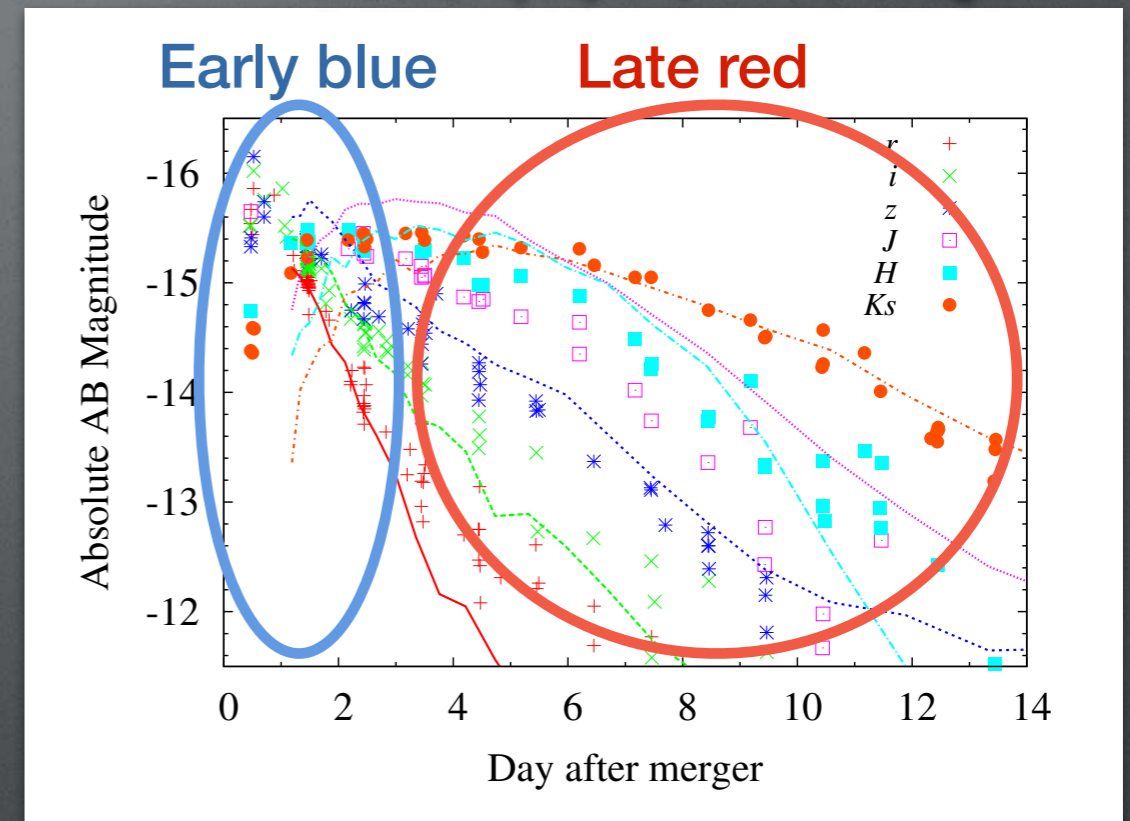
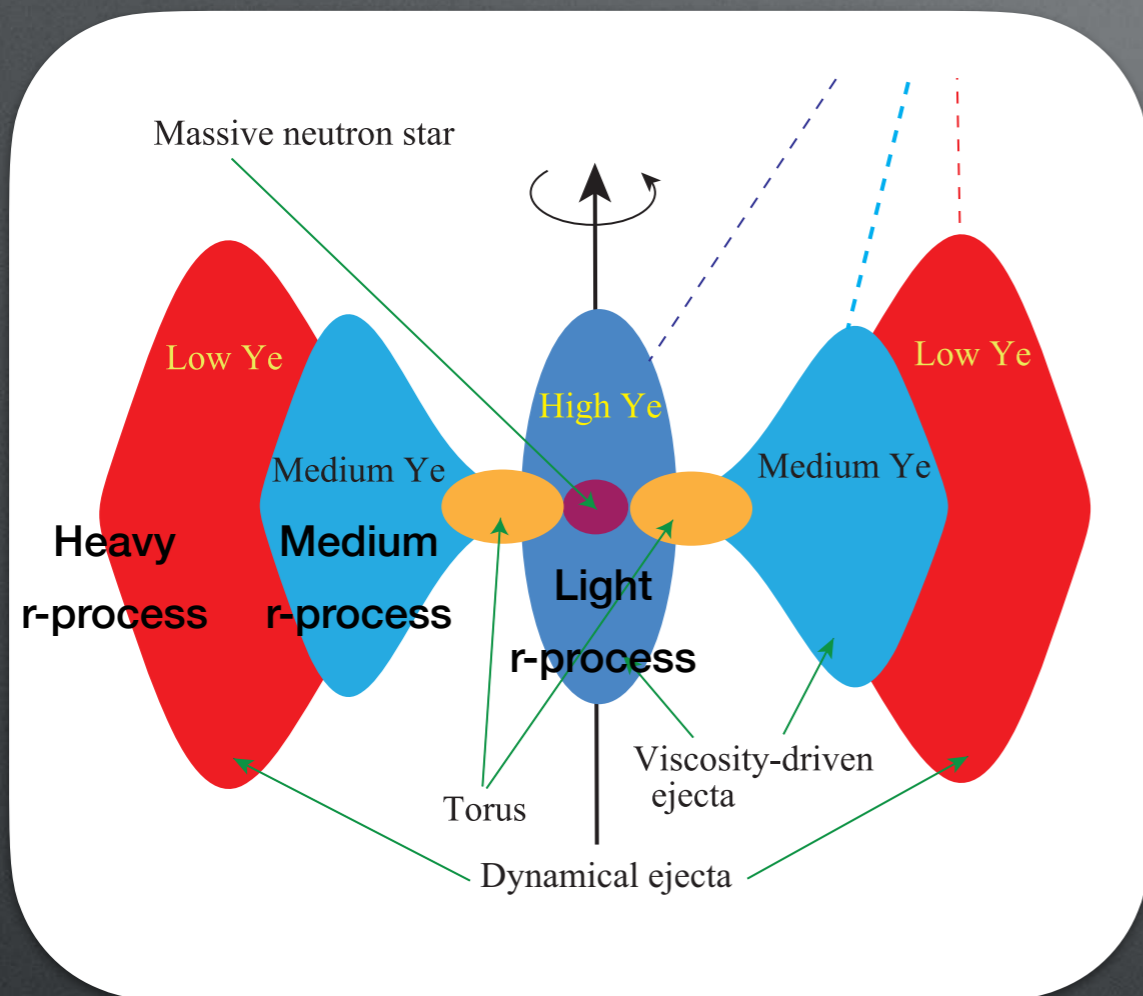


- The photosphere is in the middle ( $v \sim 0.1 - 0.2c$ ).
- The photospheric temperature is  $T \sim 10^3 \text{ K}$ .
- The spectrum peaks at optical-nIR. Rich observed data



# Kilonova Radiation Transfer Result

Shibata, Fujibayashi, KH+17, Kawaguchi+19

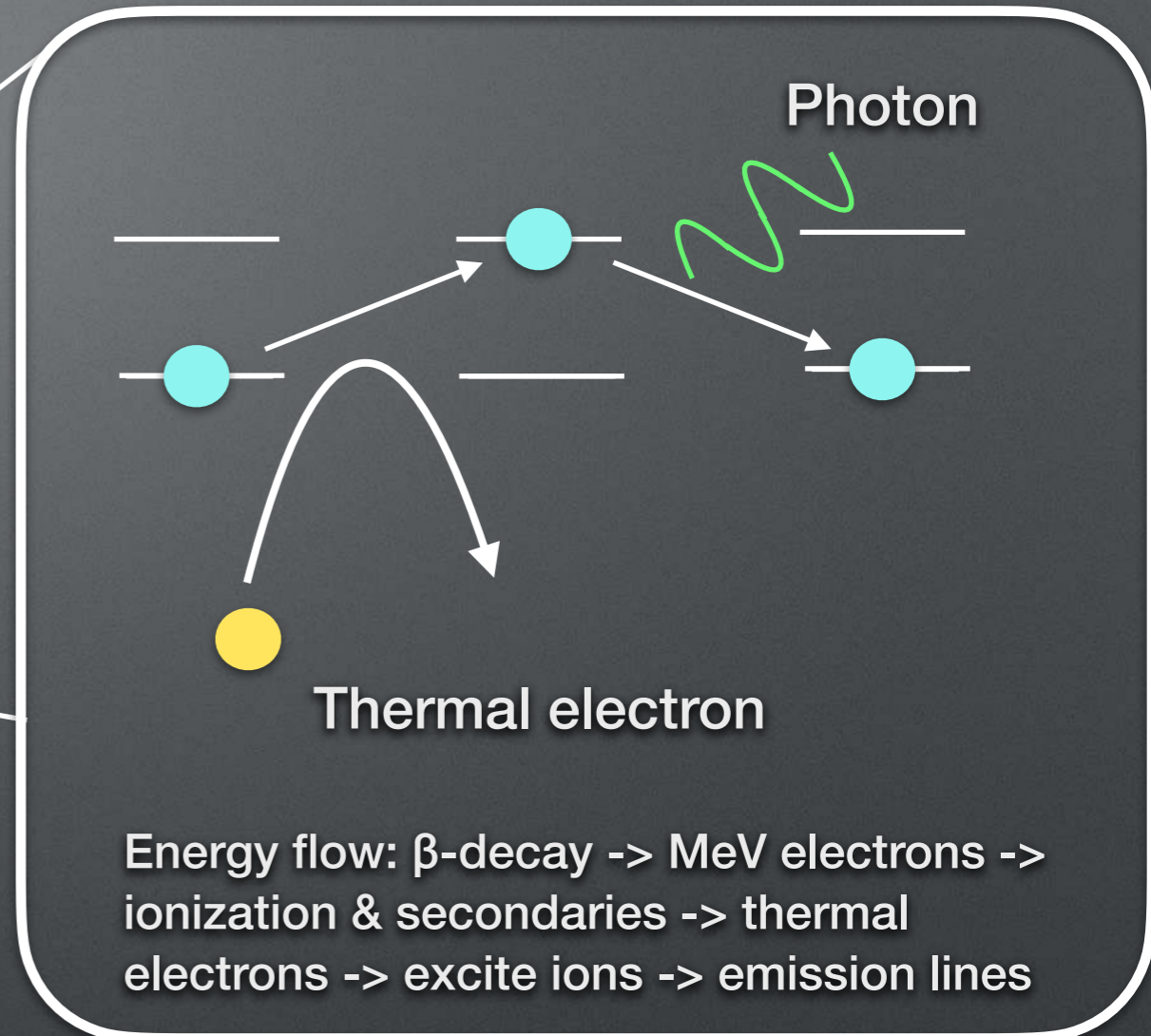
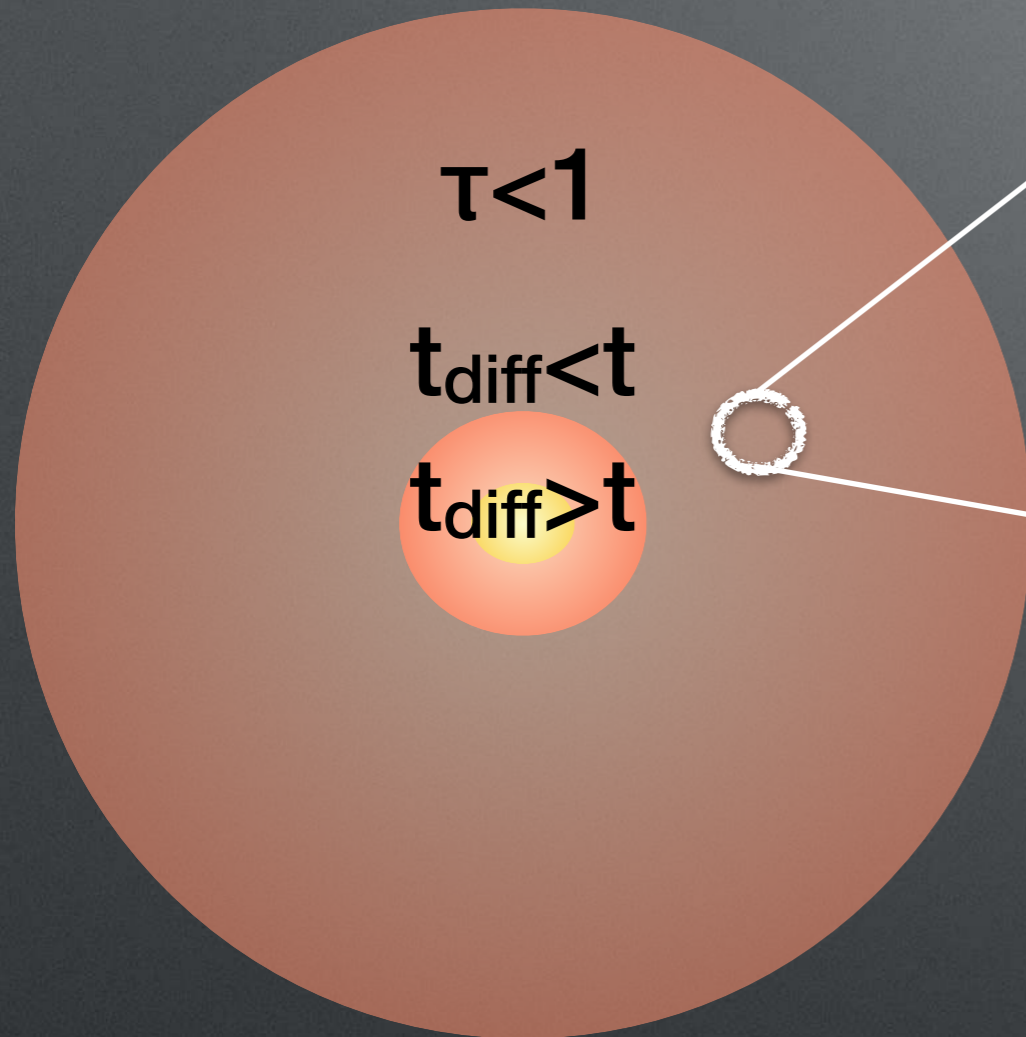


- From the observed data, the opacity must increase with time.
- The **early blue emission** ( $\sim 1$  day)  $\Rightarrow$  low opacity, e.g., **lanthanide-free** light r-process.
- The **late red emission** ( $\sim 5$  days)  $\Rightarrow$  high opacity, e.g., **lanthanide-rich** heavy r-process.
- However, the models are generally too red. We probably miss something.
- We'll hear more about the kilonova spectrum from Domoto-san.



# Kilonova Nebular Phase

2) Nebular phase  $t > 10$  day)

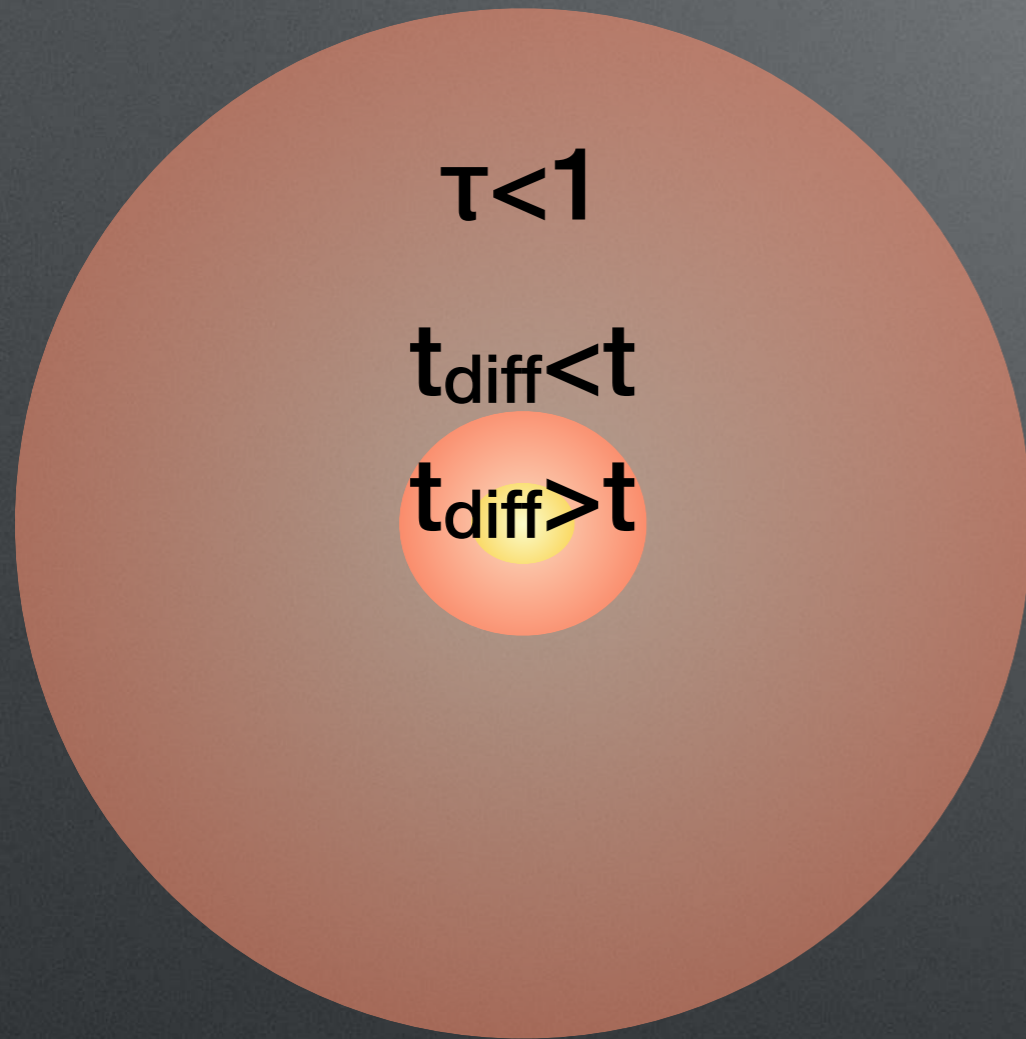


- Photons escape from the almost entire ejecta without absorption, i.e.,  $\tau < 1$ .
- Kilonova radiation is dominated by emission lines, which are narrower.
- Although it is fainter, we may have good chance to identify more elements.

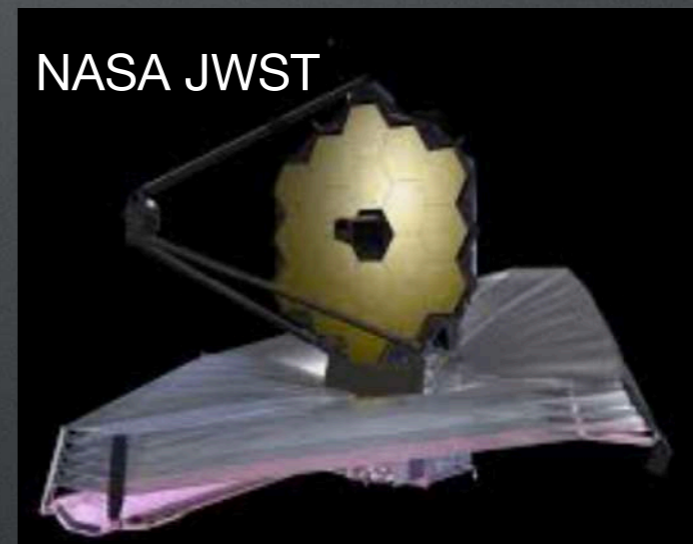
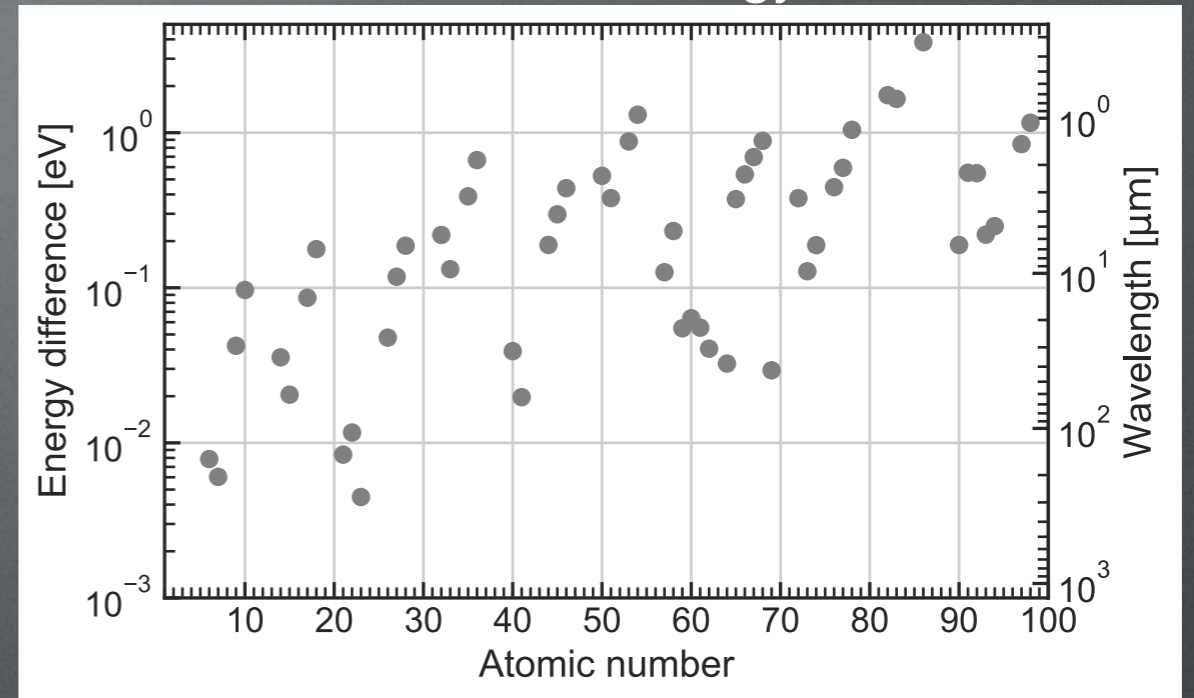


# Kilonova Nebular Phase

2) Nebular phase  $t > 10$  day)



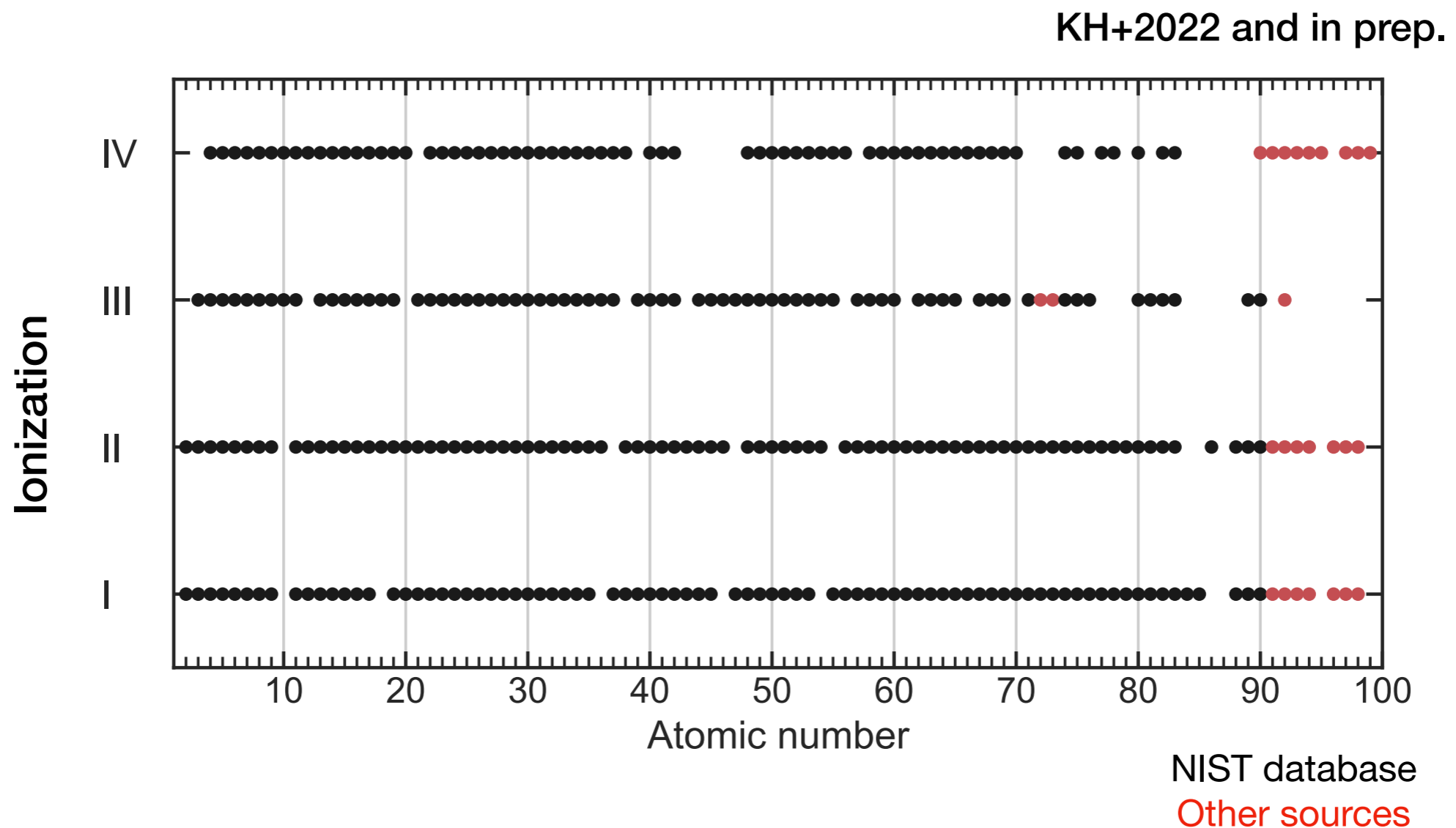
Fine Structure Energy Scale



- Photons escape from the almost entire ejecta without absorption, i.e.,  $\tau < 1$ .
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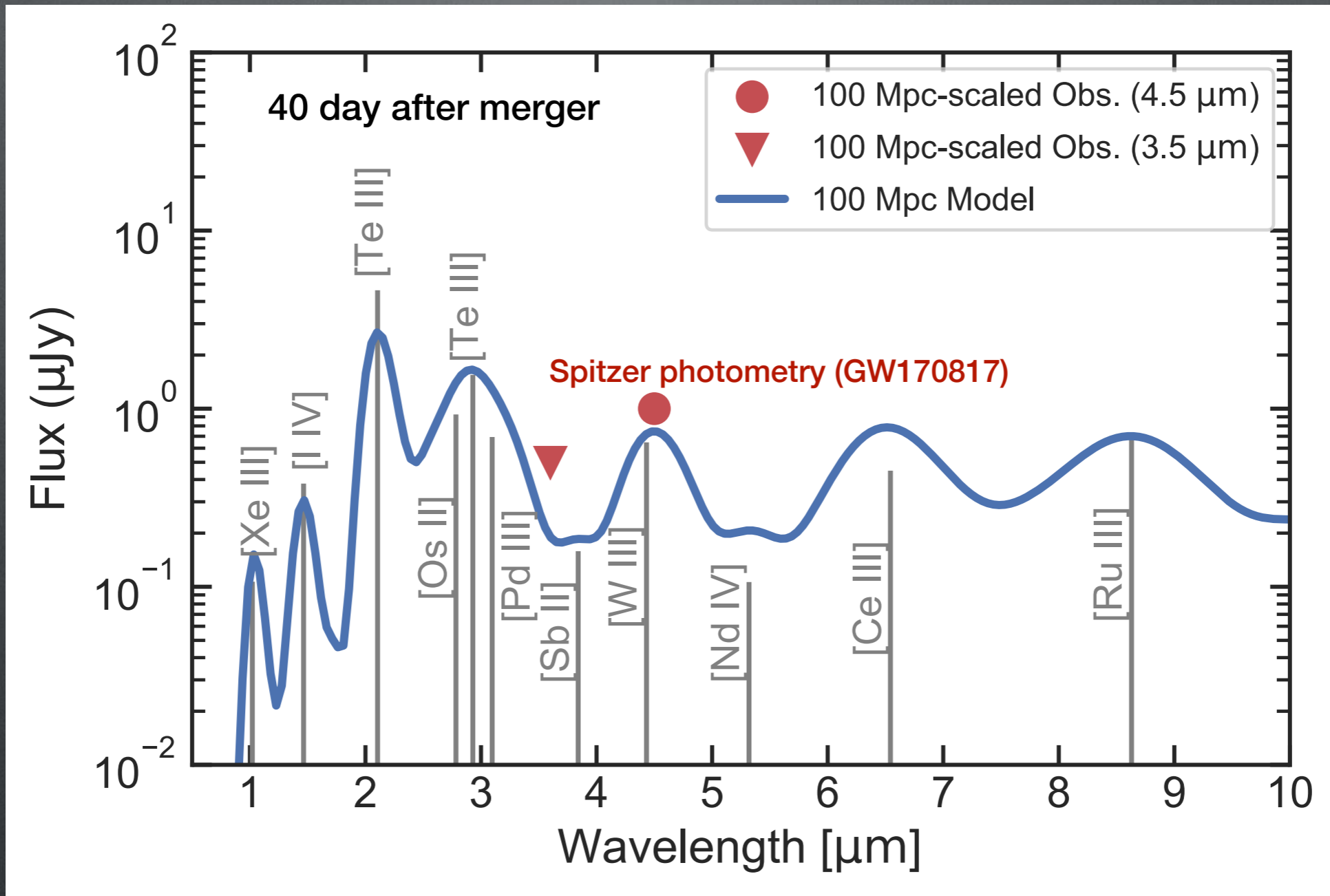
# Line list for kilonova nebula

- We constructed a forbidden (M1) line list up to Einsteinium ( $Z=99$ ).
- The experimentally calibrated levels and the LS selection rules are used.
- $A$  values are assigned with an analytic formula (Pasternack 40, Shortley 40, Bahcall & Wolf 68)
- Some ions are missing because the energy levels are unknown.





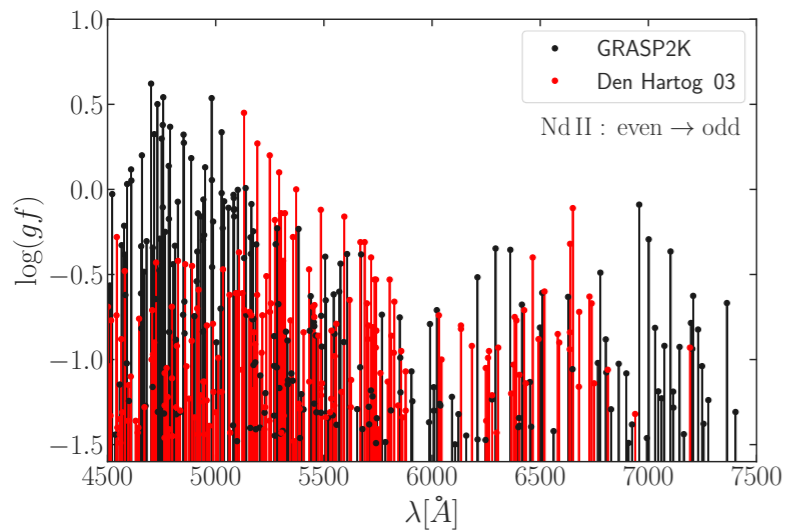
# Kilonova nebular spectrum: 40 days



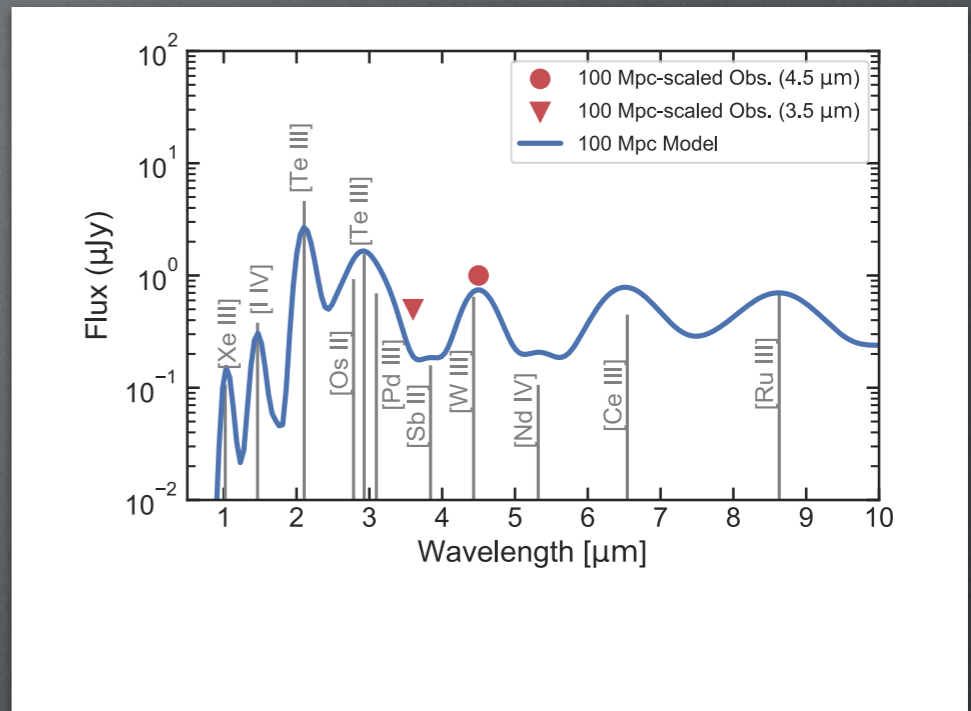
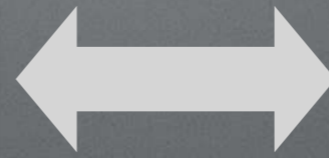
- We can now generate synthetic nebular spectra with the accurate line location.
- JWST can resolve the emission lines for kilonovae at  $\sim 100\text{Mpc}$ .
- With this model and data, the amounts of ions can be estimated.

# Forward & Reverse modelings

## Forward model

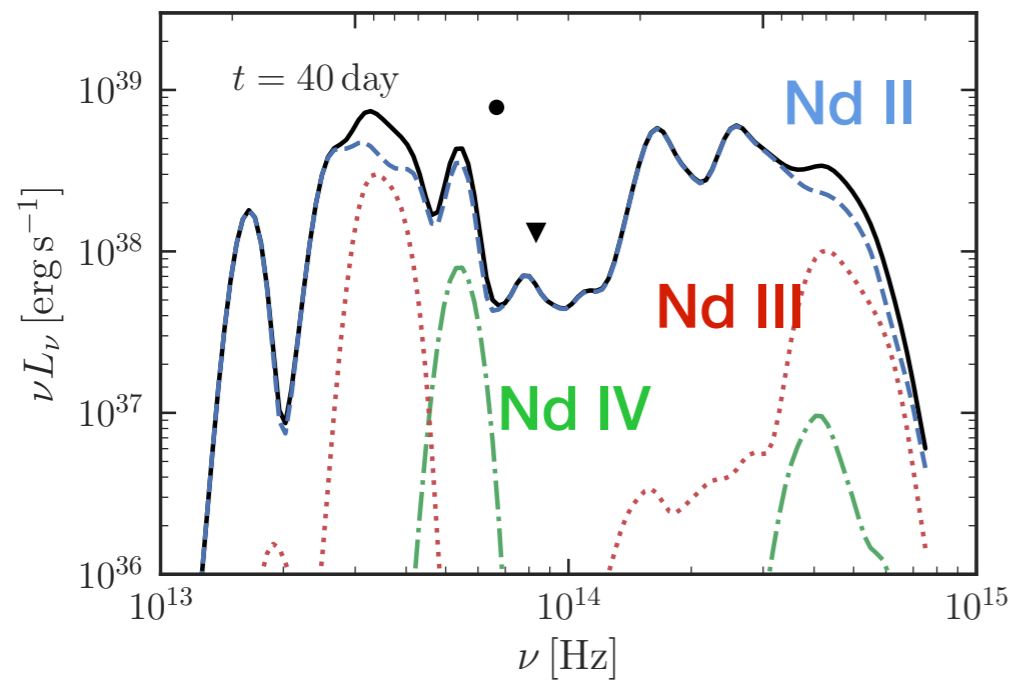
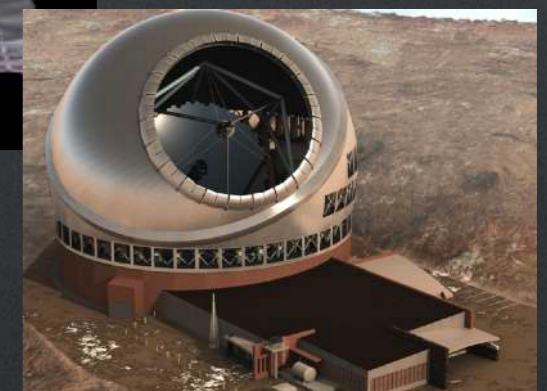
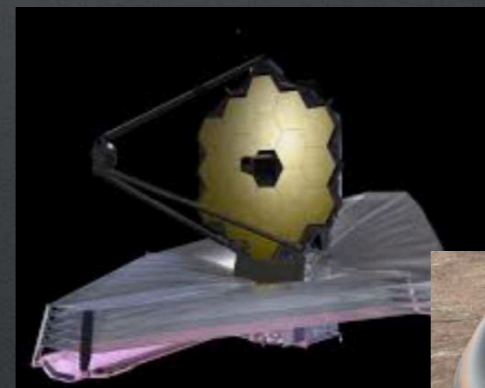


Starting from atomic data



Synthetic model

Observation



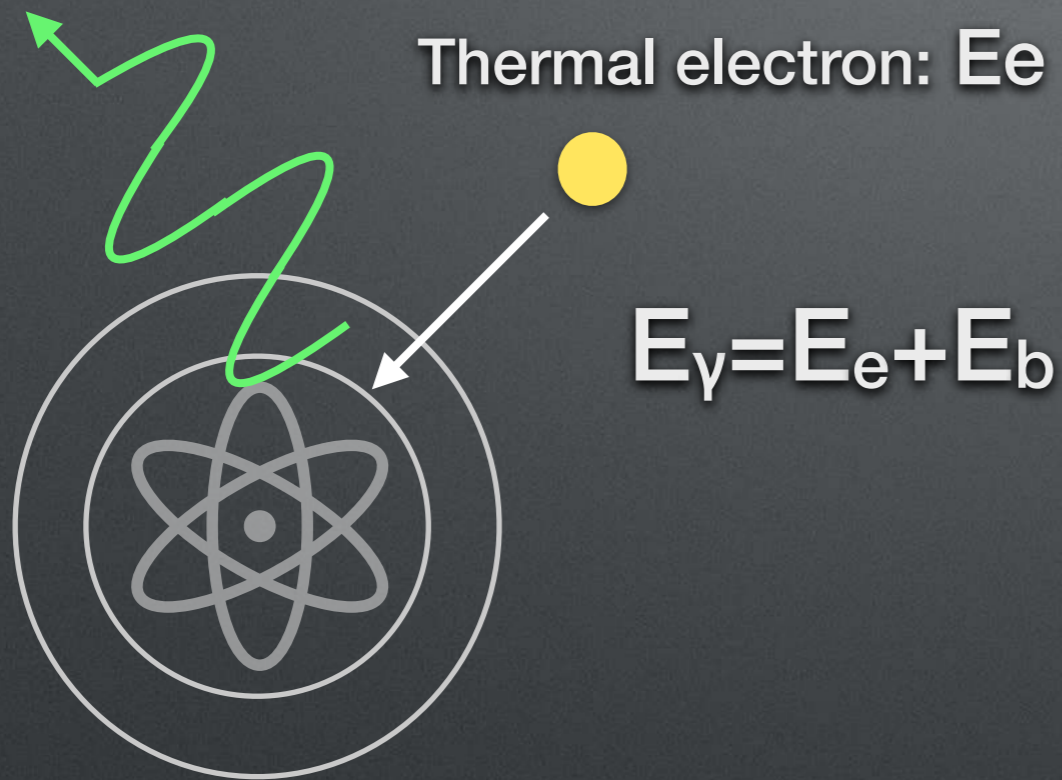
Forward modeling (KH+21)

Currently, the model is limited to a few elements.

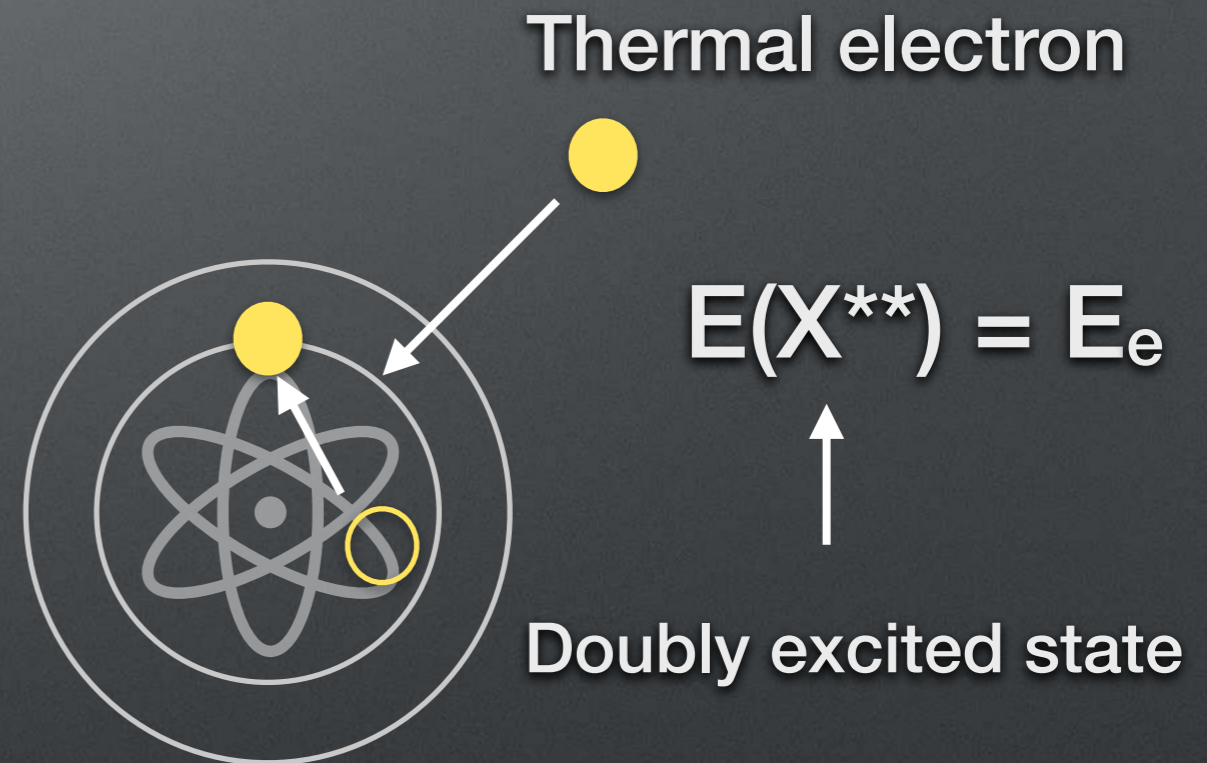


# 再結合係数の問題

## 1, 放射再結合



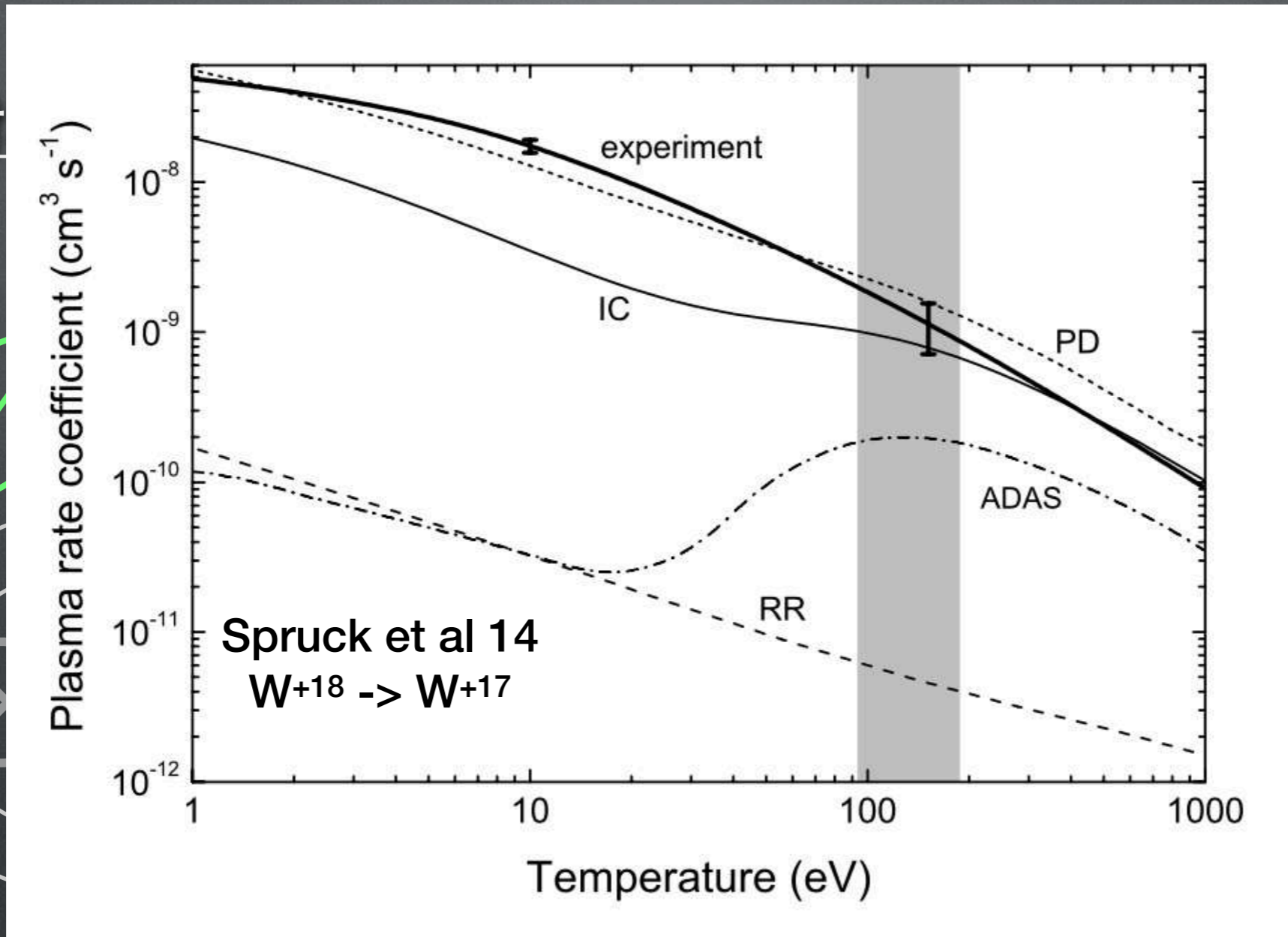
## 2, 二重電子再結合 (resonant process)





# 再結合係数の問題

1, 放射



再結合  
(process)

thermal electron

$$E(X^{**}) = E_e$$

poly excited state

- 重元素はしばしば二重電子過程が放射過程よりも圧倒的に大きい。
- 原子コードで計算可能だが、キロノバの温度(<1eV)では信頼性が低い。



# Summary

- Mass ejection of  $O(0.01 M_{\text{sun}})$  is expected from numerical relativity.
- A kilonova is powered by radioactivity of neutron-rich nuclei.
- The observed kilonova light curve points to many radioactive species of  $0.05 M_{\text{sun}}$ .
- The late-time Spitzer observation is indicative of the existence of W.
- JWST will cover a good range of the kilonova nebular spectrum.